Strongly anisotropic turbulence, structure and cascade Stratified and rotating turbulence, mainly

Summer School, Cargèse

Claude Cambon (with our team "ondes et turbulence") Laboratoire de Mécanique des Fluides et d'Acoustique CNRS – ECL – UCB Lyon I

August 15 2007

CTRL-L switch

Collaboration and documents

- Team : Fabien S. Godeferd, Julian F. Scott, Lukas Liechtenstein, Alex Delache, Guillaume Simon, Benjamin Favier
- Future or in progress : Wouter Bos, (Luminita Danaila, Jan-Bert Flor, ANISO ?) Jose Redondo, Peter A. Davidson ...
- Documents (in addition to, or instead of, published articles)
 -) http://www.lmfa.ec-lyon.fr/Fabien.Godeferd/perso/ (Summer School in Barcelone)
 -) http://www.lmfa.ec-lyon.fr/Henri.Benard/news/ (Sagaut/Cambon book)

General considerations : what is the cascade ?



not at all a dichotomic splitting, as in Frisch's book, but instead a complex effect of stretching, folding, engulfment, enrollment of velocity gradients, plus realistic vortex breakdown, reconnection, pairing ...

Quantifying the cascade

 $\bullet\,$ 'Modern' approach : Kolmogorov 4/5 law

$$\langle (\delta u_L)^3 \rangle = -\frac{4}{5}\varepsilon r + 6\nu \frac{\partial}{\partial r} \langle (\delta u_L)^2 \rangle$$

 $R_{LLL}(r) = \langle (\delta u_L)^3 \rangle$ is the key (third-order two-point) quantity

• 'Old fashion' approach : Lin equation

$$\left(\frac{\partial}{\partial t} + 2\nu k^2\right)E = T$$

T(k) can be related to the exact dynamics of *triads*, third-order three-points

Dependency on energy injection

- Realistic way: instabilities at large scale, ANISOTROPIC and inhomogeneous process,
- alteration of cascade by phase-mixing of waves, coexistence of strong and weak (wave) turbulence : HAT
- Artificial forcing, generally narrow-band in Fourier space : realistic only in strict HIT (e.g. good results from highest resolution DNS by Kaneda and coworkers)



Production ? vs. waves and neutral modes

 Homogeneous turbulence in the presence of a mean (base) flow with uniform mean gradients

$$U_i = A_{ij}(t)x_j + u_i, \qquad A_{ij} = S_{ij} + \epsilon_{ijn}W_j$$

"RDT" (Batchelor, etc), Craya (1958), Rogallo (1981), come-back to linear stability, Craik, Bayly (1986), etc.

- "production" by S : linear instability (eliptical and hyperbolical), possible *nonlinear* exponential growth (pure shear S = W/2). Homogeneity ??
- 'Pure' rotation $\mathbf{S}=0 \to$ projecting onto the rotating frame (Coriolis), inertial wave turbulence
- Similar cases without production : stable stratification with and without rotation, MHD with external magnetic field. Strong anisotropy (axisym) with homogeneity at all scales.



- Strong anisotropy at all scales, waves.
- Low Rossby (rapidly) rotating turbulence : inertial wave turbulence
- Stably stratified turbulence : toroidal cascade
- Others, MHD, weakly compressible HIT

Strong anisotropy

Anisotropic description

- ANISOTROPY/ inhomogeneity/ Intermitency
- structure functions or correlations, two-point : $R_{ij}(\mathbf{r}) = \langle u_i(\mathbf{x}) u_j(\mathbf{x} + \mathbf{r}) \rangle$



-) Single-point: componentality only

- -) Two-point: directional anisotropy
- Low dimension parameterization, SO(3) symmetry group (Arad et al., PRE, 1999)

Anisotropic description. 3D Fourier space

Anisotropic scalar (e. g. spherical harmonics) for both 'physical' and 'spectral'

$$\frac{1}{2}R_{ii}(\boldsymbol{r}) \rightarrow \frac{1}{2}\hat{R}_{ii}(\boldsymbol{k}) = e(\boldsymbol{k})$$

$$\sum r_n^m(r)Y_n^m(\theta_r,\phi_r) \rightarrow \sum \varphi_n^m(k)Y_n^m(\theta_k,\phi_k)$$

Avoiding a 'schizophrenic' viewpoint ! (Cambon & Teissèdre 1985, CRAS Paris)

• A trace-deviator decomposition restricted to solenoidal space

$$\hat{R}_{ij} = \underbrace{U(k)P_{ij}}_{\text{isotropic}} + \underbrace{\mathcal{E}(k)P_{ij}}_{\text{directional}} + \underbrace{\Re\left(Z(k)N_iN_j\right)}_{\text{polarization}}.$$

(Cambon & Jacquin, JFM, 1989), $P_{ij} = \delta_{ij} - rac{k_i k_i}{k^2}$, $m{N}$ 'helical mode'. Helicity ?

Summer School, Cargèse

August 15 2007

Claude Cambon (with our team "ondes et turbulence")



Rotating turbulence, MHD simplified case, $m{k} \perp \hat{m{u}}$

Summer School, Cargèse

Stable stratification and rotation

Incompressible N-S with Coriolis + buoyancy $m{F}=2m{\Omega} imesm{u}+bm{n}$ $m{n}$ local upstream vertical, Coriolis $fm{n} imesm{u}$

 $\dot{b} + N^2 u_{\parallel} - P_r
u
abla^2 b = 0$ (idem passive scalar with mean scalar gradient)

Cartoon of oscillations, towards waves



Two external parameters

(frequencies) : N and f.

Summer School, Cargèse

Eigenmodes decomposition

• incompressibility and pressure \rightarrow 3D Fourier space $(\boldsymbol{u},b)(\boldsymbol{x},t)=$

$$\sum e^{\imath \boldsymbol{k} \cdot \boldsymbol{x}} \left(\underbrace{a_{0} \boldsymbol{N}^{(0)}}_{\text{vortex (QG)}} + \underbrace{a_{+1} \boldsymbol{N}^{(1)} e^{\imath \sigma_{k} t} + a_{-1} \boldsymbol{N}^{(-1)} e^{-\imath \sigma_{k} t}}_{\text{wave (AG)}} \right)$$

- Dispersion law $\sigma_k = \sqrt{N^2 \sin^2 \theta + f^2 \cos^2 \theta}$
- Linear dynamics: slow amplitudes $a_{0,\pm 1}$ are constant. Nonlinear case
- Advantages $\mathbf{k} \cdot \hat{\mathbf{u}} = 0$, five $(u_1, u_2, u_3, b, p) \rightarrow$ three (a_0, a_{+1}, a_{-1}) .

ref. Cambon et al., Bartello, Smith & Waleffe, Morinishi, Kaneda ... etc

Wave aspects



Mc Ewan (1967), Mowbray &

Rarity (1967), Godeferd & Lollini, JFM (1999)

Summer School, Cargèse

Rotation, phase mixing: Linear and nonlinear

- Relevance of linear solution depends on the order and type of statistical correlations
 - -) doubles: 2 point 1 time: $e^{\imath \sigma_k t}$, $e^{-\imath \sigma_k t}$
 - -) doubles: 2 point 2 time: $e^{i\sigma_k(t\pm t')}$
 - -) triples: 3 point 1 time: $e^{i(\pm \sigma_k \pm \sigma_p \pm \sigma_q)t} \rightarrow \text{nonlinear} \cdots$

$$\langle \omega_3^3 \rangle = \sum \int \exp[ift(\cos\theta_k + s'\cos\theta_p + s''\cos\theta_q] S(\boldsymbol{k}, \boldsymbol{p}, \epsilon t) d^3 \boldsymbol{p} d^3 \boldsymbol{k}$$

($\cos \theta_k = k_3/k$) Need for initial triple correlations at THREE point. Many other correlations.



Cyclonic / anticyclonic asymmetry: (64^3 LES ?) Bartello et al. (1994), Morize et al. (2005), Gence & Frick (2001), Staplehurst et al. (2007), van Bokhoven et al. (2006). No need for centrifugal inst.

$$rac{d}{dt}\langle\omega_3^3
angle=6\Omega\langlerac{\partial u_3}{\partial x_3}\omega_3^2
angle+4$$
th-order, viscous

Summer School, Cargèse

Exact spectral equations for nonlinear theory

$$\left(\frac{\partial}{\partial t} + 2\nu k^2\right)e(k,\cos\theta,t) = T^{(e)}(k,\cos\theta,t)$$

Poincaré transformation $\hat{u}_i(\mathbf{k}, t) = \sum_{s=\pm 1} N_i(s\mathbf{k}) e^{isft\cos\theta} a_s(\mathbf{k}, \epsilon t)$,

$$T^{(e)} = \int_{t_0}^t \sum_{s,s',s''=\pm 1} \int_{p+q=k} e^{if(t-t')(s\cos\theta_k + s'\cos\theta_p + s''\cos\theta_q)}$$

$$S_{ss's''}(\boldsymbol{k},\boldsymbol{p},t,\epsilon t')d^{3}\boldsymbol{p}dt',$$

with exact tradic conservation $\dot{a}_s(\mathbf{k}) = (s'p - s''q)G_{kpq}a^*_{s'}(\mathbf{p})a^*_{s''}(\mathbf{q})$, possibly restricted to resonant triads.

August 15 2007 Claude Cambon (with our team "ondes et turbulence")

20

Summer School, Cargèse

Results. NONLINEAR statistical theory

• From classical EDQNM (isotropic, no rotation, Orszag 1970, Bos & Bertoglio, 2006)



• ... to EDQNM3 \rightarrow (A) QNM energy equation (Bellet *et al.*, JFM, 2006)

Summer School, Cargèse

Angle-dependent spectrum

• Isotropy breaking by spectral transfer $T^{(e)}(\mathbf{k})$: directional anisotropy:



• Spherical averaging $\rightarrow E(k, t_f)$, prefactor $E \sim \frac{\Omega}{t} k^{-3}$, not 2D !

Summer School, Cargèse

AQNM and DNS



 512^3 DNS by Liechtenstein *et al.*, JOT, 2005

Summer School, Cargèse

Inertial wave-turbulence, resonant interactions, 2D or not 2D

 Low dimension of active manifolds : overestimated in forced ? DNS/LES ? 'TRUE' 2D embedded in 3D : a DIRAC singularity

$$E(k) \sim f^2 k^{-3}, \qquad e(k_{\perp}, k_{\parallel}) = \frac{E(k_{\perp})}{2\pi k_{\perp}} \delta(k_{\parallel})$$

• integral singularity from theoretical wave-turbulence

$$\begin{split} e(k_{\perp}k_{\parallel}) \sim k_{\parallel}^{-1/2} k_{\perp}^{-7/2} &= k^{-4} x^{-1/2} \quad \text{Galtier 2002} \\ e(k_{\perp}k_{\parallel}) \sim k_{0}^{-1/2} k_{\parallel}^{-1/2} k_{\perp}^{-3} &= k_{0}^{-1/2} k^{-7/2} x^{-1/2} \quad \text{CRG 2004} \\ E(k) \sim \frac{f}{t} k^{-3}, \quad e(k, x \sim 0) \sim k^{-4} \quad \text{BGSC-2006} \end{split}$$

Summer School, Cargèse

Reintroducing stratification : ANTI 2D and toroidal (strong) cascade

Summer School, Cargèse

Geometric aspects of eigenmodes



Summer School, Cargèse



Summer School, Cargèse



Angle-dependent toroidal and poloidal modes (Liechtenstein, 2006)

Summer School, Cargèse



Summer School, Cargèse

Pure stratification: scaling arguments

- 2D or not 2D ? Charney (1971), Lilly (1983), Lindborg (1999) using third order structure functions from observations.
- Froude numbers, horizontal and vertical, $Fr_h = U/(NL_h)$, $Fr_v = U/(NL_v)$, $L_h \gg L_v$.
- $L_v \sim U/N$ ("zig-zag" instability ? Billant & Chomaz 1999, 2002) $\rightarrow Fr_v \sim 1$ (to contrast with Riley et al. 1981 ?)
- Proposed scaling (Linborg 2006) $E_{hh}(k_h) = C_1 \varepsilon_k^{2/3} k_h^{-5/3}$, $E_p(k_h) \sim C_2 \epsilon_p \epsilon_K^{-1/3} k_h^{-5/3}$, $E_{vv} \sim N^2 k_v^{-3}$.
- $Ri \sim 1/4$ threshold vs. Fr^2Re scaling (Riley) ?, rediscovery of polo-toro decomposition (Brethouwer, Linborg, Billant, Chomaz.)

Pure stratification: Generalized Lin equations

$$\left(\frac{\partial}{\partial t} + 2\nu k^2\right) e^{(tor)} = T^{(tor)} \tag{1}$$

$$\left(\frac{\partial}{\partial t} + 2\nu k^2\right)e^{(w)} = T^{(w)} \tag{2}$$

$$\left(\frac{\partial}{\partial t} + 2\nu k^2 + 2iN\frac{k_\perp}{k}\right)Z' = T^{(z')} \tag{3}$$

Energy spectra $e^{(tor),(pol),(pot)}(k_{\perp},k_{\parallel})$, imbalance deviator Z', $\Re Z' = e^{(pol)} - e^{(pot)}, e^{(w)} = e^{(pol)} + e^{(pot)}$.

A lot of information can be generated, vs. second and third-order structure functions.

Summer School, Cargèse

The toroidal cascade. Why not 2D?

- Why the toroidal component only ? $e^{(1)} \cdot \dots \frac{\partial u}{\partial t} + \omega \times u + \nabla \left(p + \frac{u^2}{2} \right) bn$, $\dot{u}^{(1)} + e^{(1)} \cdot \sum_{p+q=k} (\hat{\omega}(p) \times \hat{u}(q)) = 0$, $\hat{u} = u^{(1)} e^{(1)} + u^{(2)} e^{(2)}$, $\hat{\omega} = \imath k \left(u^{(1)} e^{(2)} - u^{(2)} e^{(1)} \right)$
- Following Kraichnan and Waleffe (1992, 1993): stability of a single triad

$$\dot{u}_k^{(1)} = (p_\perp^2 - q_\perp^2) G u_p^{(1)*} u_q^{(1)*}, \tag{4}$$

$$\dot{u}_{p}^{(1)} = (q_{\perp}^{2} - k_{\perp}^{2}) G u_{q}^{(1)*} u_{k}^{(1)*},$$
(5)

$$\dot{u}_q^{(1)} = (k_\perp^2 - p_\perp^2) G u_k^{(1)*} u_p^{(1)*}, \tag{6}$$

 quasi 2D or not, reverse or direct cascade ? cylinder to cylinder, shell to shell, angle to angle : very rich and various morphology ...

Analogy with an 'Euler problem'

• The solid in its principle axes of inertia

$$I_1 \dot{\Omega}_1 = (I_2 - I_3) \Omega_2 \Omega_3,$$
(7)

$$I_2 \dot{\Omega}_2 = (I_3 - I_1) \Omega_3 \Omega_1 \tag{8}$$

$$I_3 \dot{\Omega}_3 = (I_1 - I_2) \Omega_1 \Omega_2,$$
 (9)

• Conservations laws, rot. kin. energy $(I_1\Omega_1^2 + ...) \rightarrow$ kin. energy (triad), norm of the angular momentum $(I_1\Omega_1)^2 + ... \rightarrow$ vertical enstrophy (triad)



• Instabilities, reverse interactions only.

Summer School, Cargèse

Revisiting the QG cascade

- Detailed conservation of QG energy and potential enstrophy (linear ?) $\frac{k^2 \sigma_k^2}{N^2} \xi^{(0)} \xi^{(0)*} = k_{\perp}^2 u^{(1)} u^{(1)*} + \left(\frac{f}{N} k_{\parallel}\right)^2 u^{(3)} u^{(3)*}$
- Reworking on N^{000} (Bartello 1995)

$$\dot{\xi}_{k}^{(0)} = (p^{2}\sigma_{p}^{2} - q^{2}\sigma_{q}^{2})G'\xi_{p}^{(0)*}\xi_{q}^{(0)*}, \tag{10}$$

$$\xi_p^{(0)} = (q^2 \sigma_q^2 - k^2 \sigma_k^2) G' \xi_q^{(0)*} \xi_k^{(0)*},$$
(11)

$$\dot{\xi}_q^{(0)} = (k^2 \sigma_k^2 - p^2 \sigma_p^2) G' \xi_k^{(0)*} \xi_p^{(0)*},$$
(12)

• 'linear' (quadratic) limit ? Ertel theorem, flat isopycnes. Dual cascade or not, why ?

Toroidal main energy drain, 'wave-released' EDQNM2



Summer School, Cargèse



August 15 2007 Claude Cambon (with our team "ondes et turbulence")

Summer School, Cargèse

Coexistence of weak and strong turbulence: MHD



strong turbulence with additional Joule dissipation effect (Moffat 1967, Alboussière, etc)

Coexistence of weak and strong turbulence: Aeroacoustics



Fauchet et al. (1997).

Summer School, Cargèse

Some concluding comments, open issues

- 1. What about structure functions and 4/5 equation for strong anisotropy at all scale ? Interest for vorticity, potential vorticity too
- 2. Anisotropic cascade : $\langle (\delta u_{\parallel})^3 \rangle$ versus $T^i(k_{\parallel},k_{\perp})$?
 - -) Exact Lin-type equations

-) A systematic way to construct triadic correlations, consistently with detailed conservation laws

- -) Need for improved ED for strong turbulence (Wouter-EDQNM, TFM, LRA ?)
- 3. Competition between waves and vortices to organise Lagrangian diffusion (also plasmas ?)
- 4. A better link between physical and spectral space ? (ANISO)

Towards a theory of axisymmetric weak+strong turbulence



- Work remain to be done for toroidal, QG, + main waves interactions (catalytic ?)
- Toro-polo with Elsasser variables, MHD, plasmas
- Not to forget DNS/LES with dedicated post-processing (Marenostrum project, etc)

Summer School, Cargèse