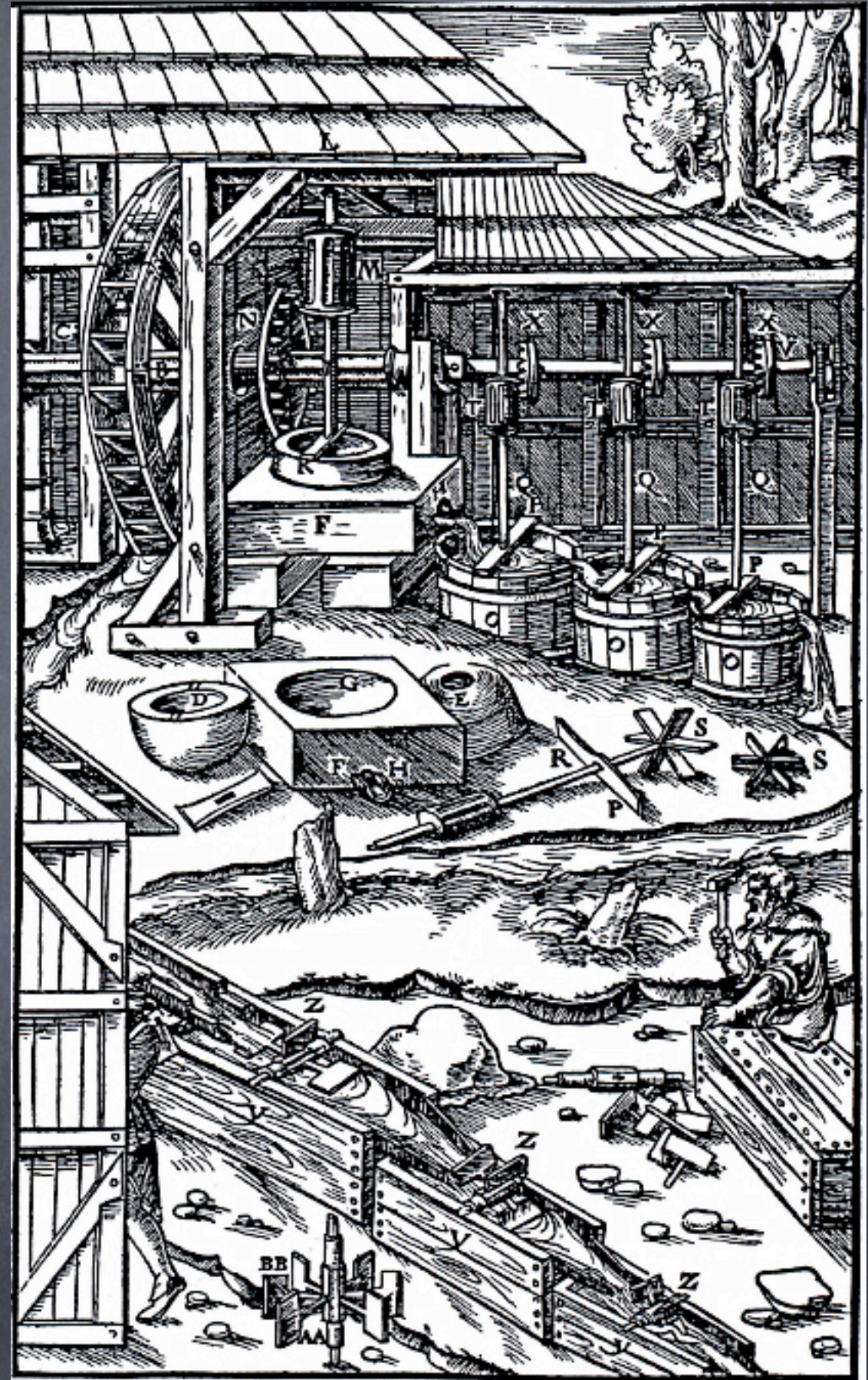
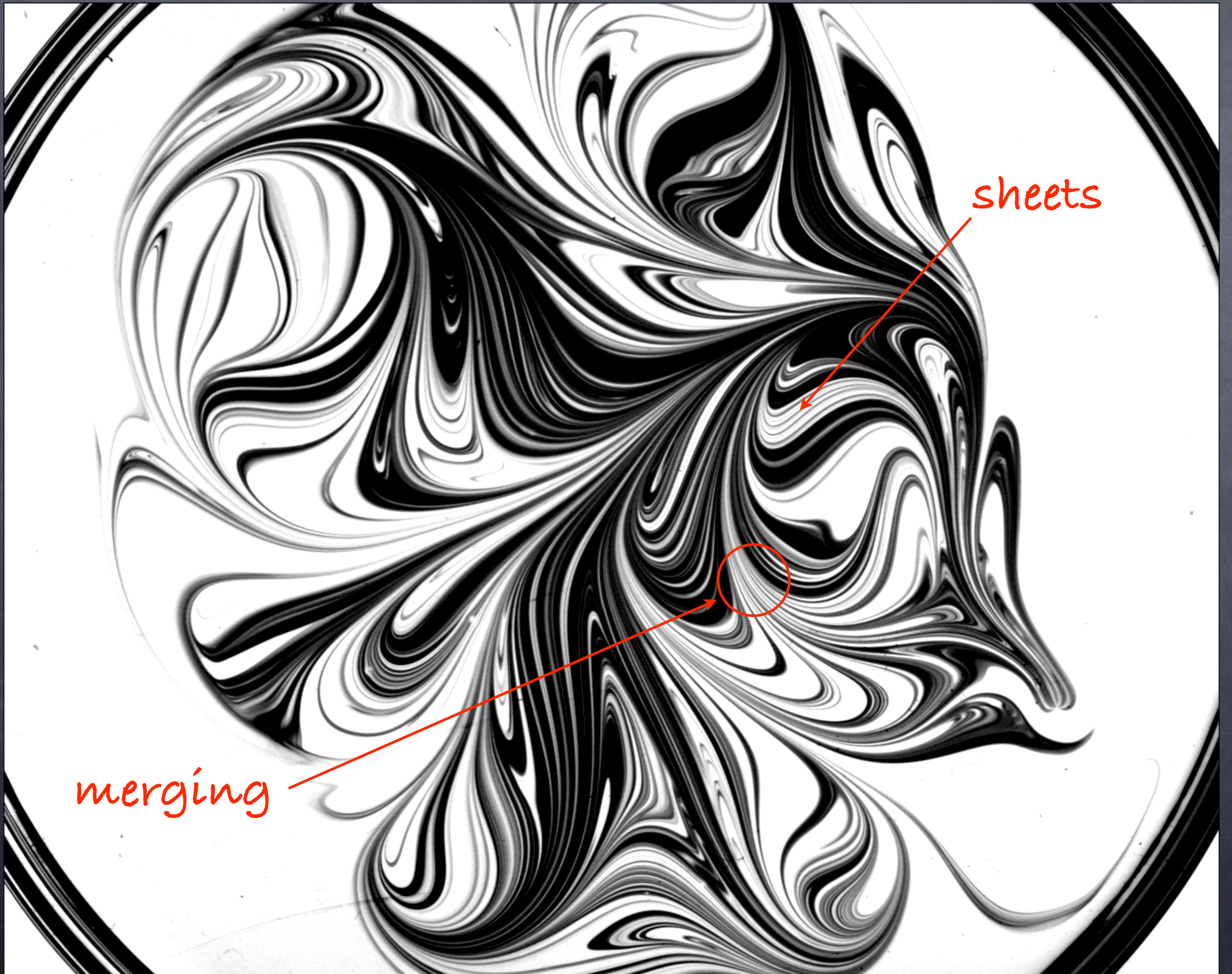


Mixing by Random
Stirring



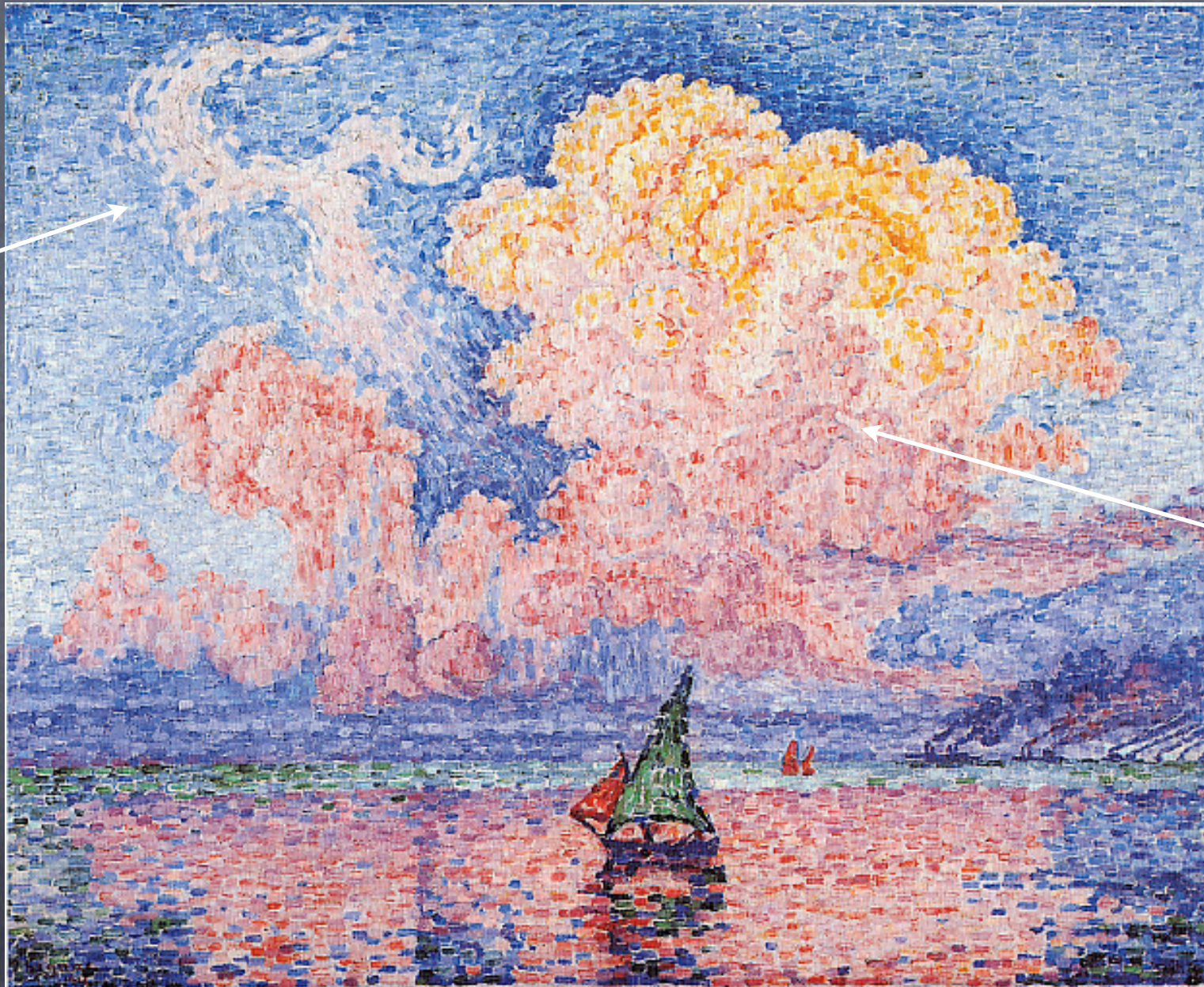


sheets

merging

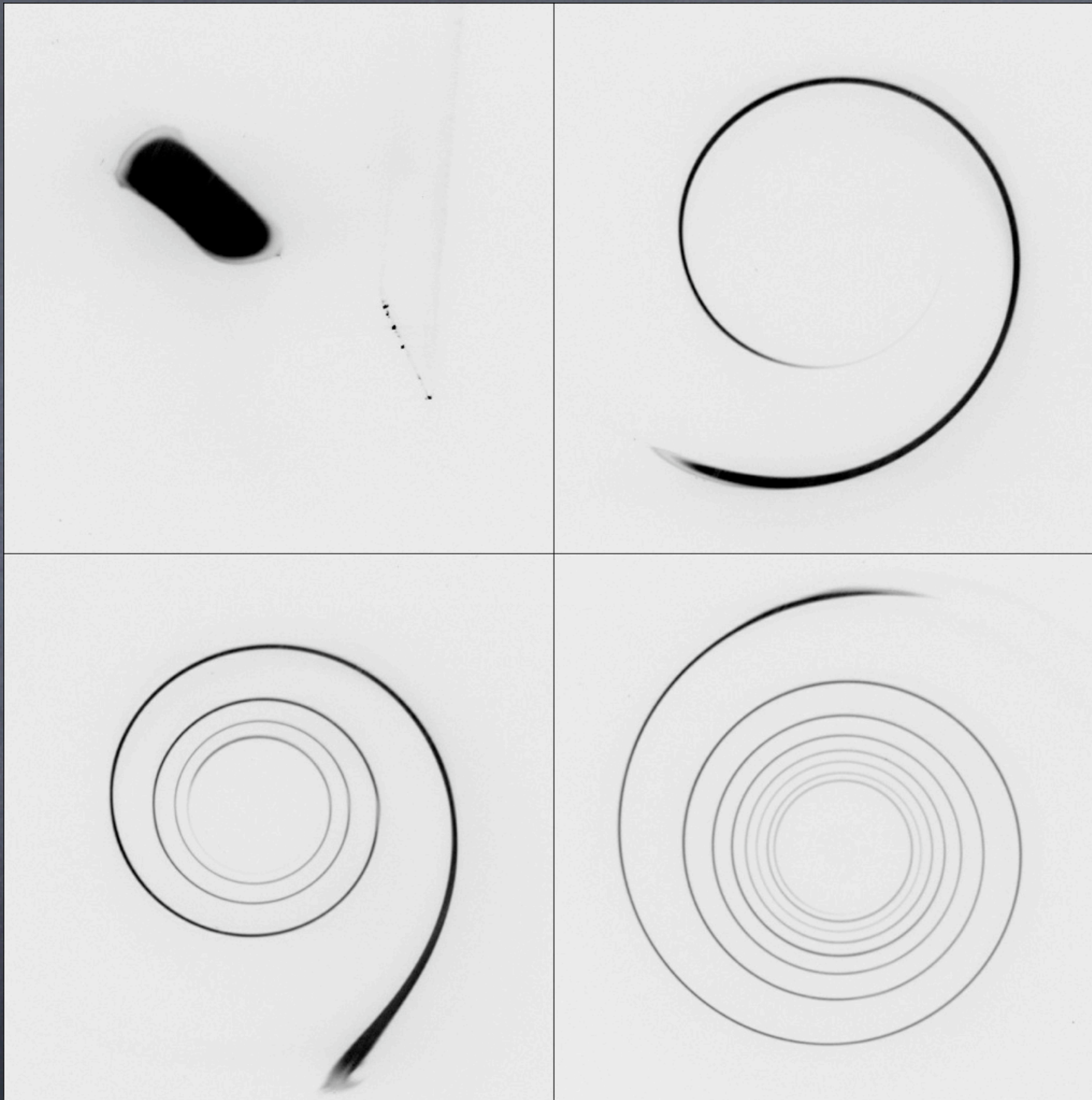
Question: $P(C) = ?$

Dispersing
Mixture



Confined
Mixture

P. Signac, 'Antibes, Le Nuage Rose' (1916)



Meunier & Villermanx, JFM 476 (2003)

Axisymmetric, diffusing vortex



$$\partial_t v_\theta = \nu \left\{ \frac{\partial^2 v_\theta}{\partial r^2} + \frac{1}{r} \frac{\partial v_\theta}{\partial r} - \frac{v_\theta}{r^2} \right\}$$

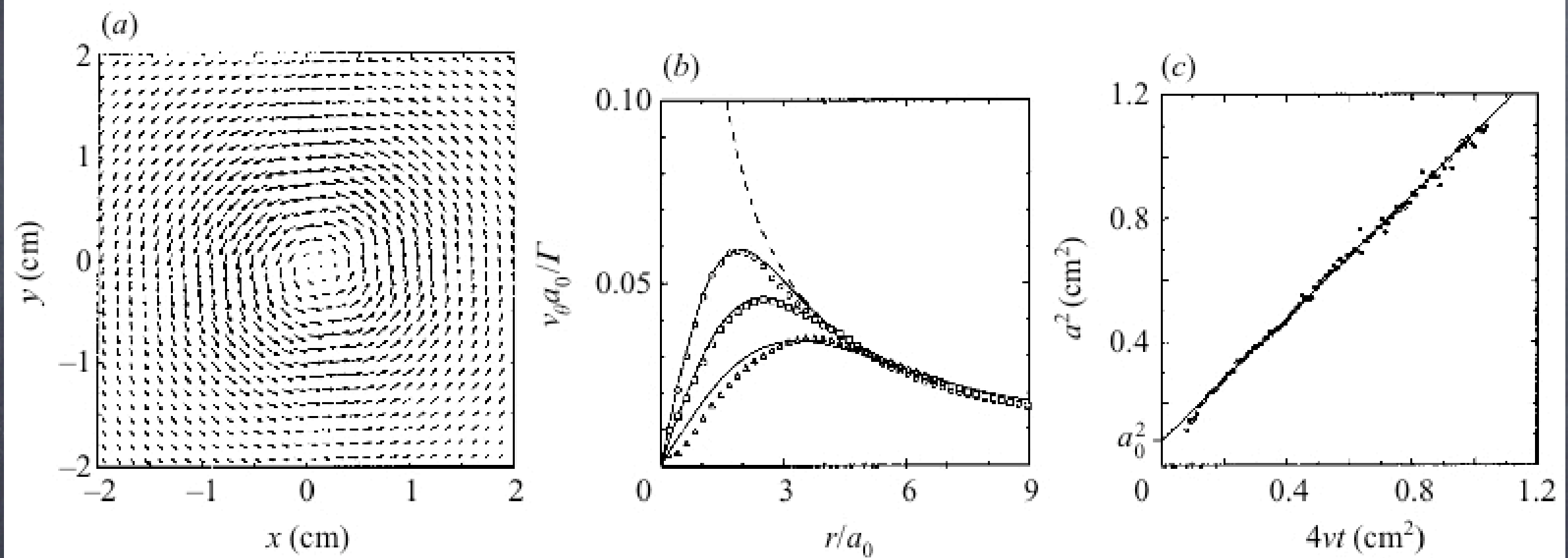
$$v_\theta = \frac{\Gamma}{2\pi r} (1 - e^{-r^2/a^2})$$

$$a^2 = a_0^2 + 4\nu t$$

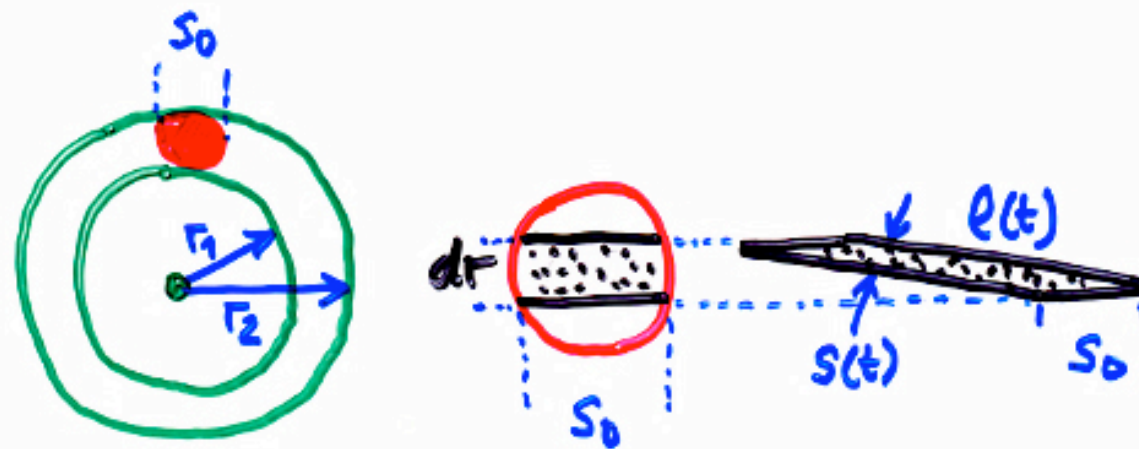
Lamb-Oseen

$$\begin{aligned} \Gamma &= 2\pi a_0 v_\theta(a_0, t=0) \\ &= \text{circulation} \end{aligned}$$

Lamb-Oseen vortex



Scalar blob deformation



$$r_1 \gg a(t) : v_\theta \approx \frac{\Gamma}{2\pi r}$$

$$\begin{aligned} \text{length: } dl(r, t) &= \sqrt{dr^2 + (r d\theta)^2} \\ &= dr \sqrt{1 + \frac{\Gamma^2 t^2}{\pi^2 r^4}} \end{aligned}$$

$$\text{thickness: } s(t) = \frac{s_0 dr}{dl(t)} = \frac{s_0}{\sqrt{1 + \frac{\Gamma^2 t^2}{\pi^2 r^4}}}$$

Spiral total length:

$$\int_{r_1}^{r_2} dl(r, t) \sim t$$

Convection - diffusion

In the local frame of a lamellae:



substrate velocity: $u = -\frac{x}{s} \frac{ds}{dt}$
 $v = \frac{y}{s} \frac{ds}{dt}$

Diffusion equation:

$$\frac{\partial c}{\partial t} + u \frac{\partial c}{\partial x} + v \frac{\partial c}{\partial y} = D \left(\frac{\partial^2 c}{\partial x^2} + \frac{\partial^2 c}{\partial y^2} \right)$$

$$u \frac{\partial c}{\partial x} / v \frac{\partial c}{\partial y} = O\left(\frac{s}{\ell}\right); \quad \frac{\partial^2 c}{\partial x^2} / \frac{\partial^2 c}{\partial y^2} = O\left(\frac{s}{\ell}\right)^2$$

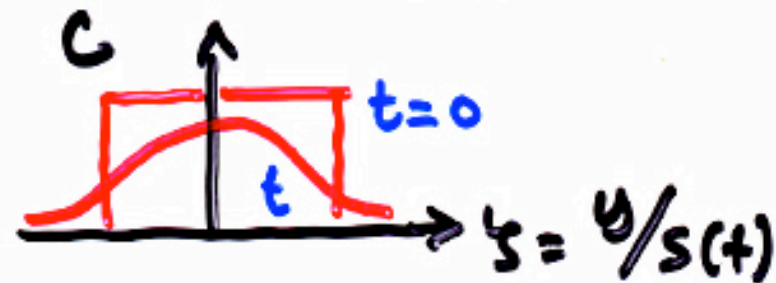
Thus: $\frac{\partial c}{\partial t} + v \frac{\partial c}{\partial y} = D \frac{\partial^2 c}{\partial y^2}$

With: $\xi = \frac{y}{s(t)}$; $\tau = D \int_0^t \frac{dt'}{s(t')^2}$

$$\boxed{\frac{\partial c}{\partial \tau} = \frac{\partial^2 c}{\partial \xi^2}}$$

Solution:

$$C(\xi, \tau) = \frac{1}{2} \left\{ \operatorname{erf} \left[\frac{\xi + 1/2}{2\sqrt{\tau}} \right] - \operatorname{erf} \left[\frac{\xi - 1/2}{2\sqrt{\tau}} \right] \right\}$$



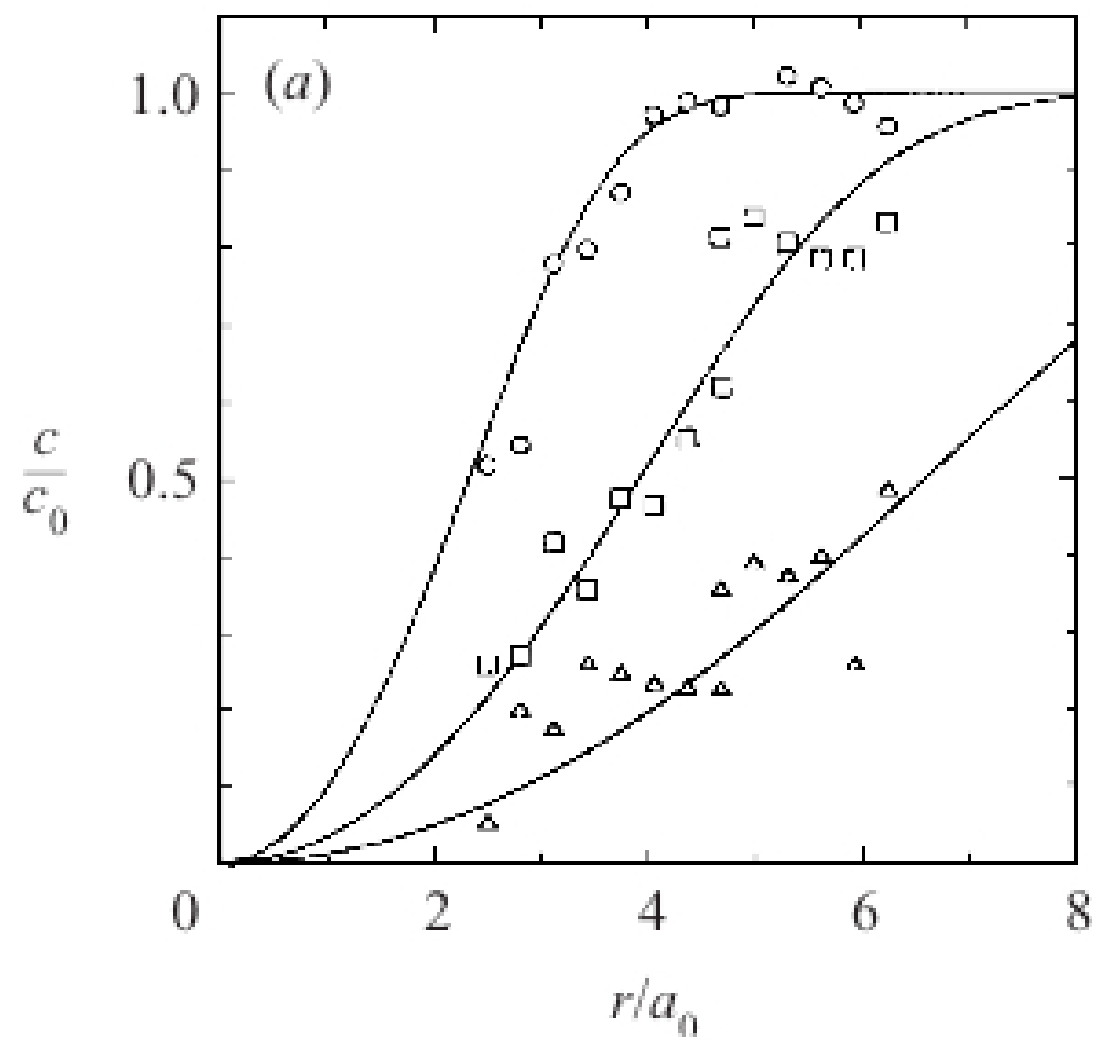
Mixing time: $t_s(r) \approx \frac{r^2}{D} \left(\frac{s_0}{r} \right)^{2/3} \left(\frac{\Gamma}{D} \right)^{1/3}$

For $t > t_s(r)$:

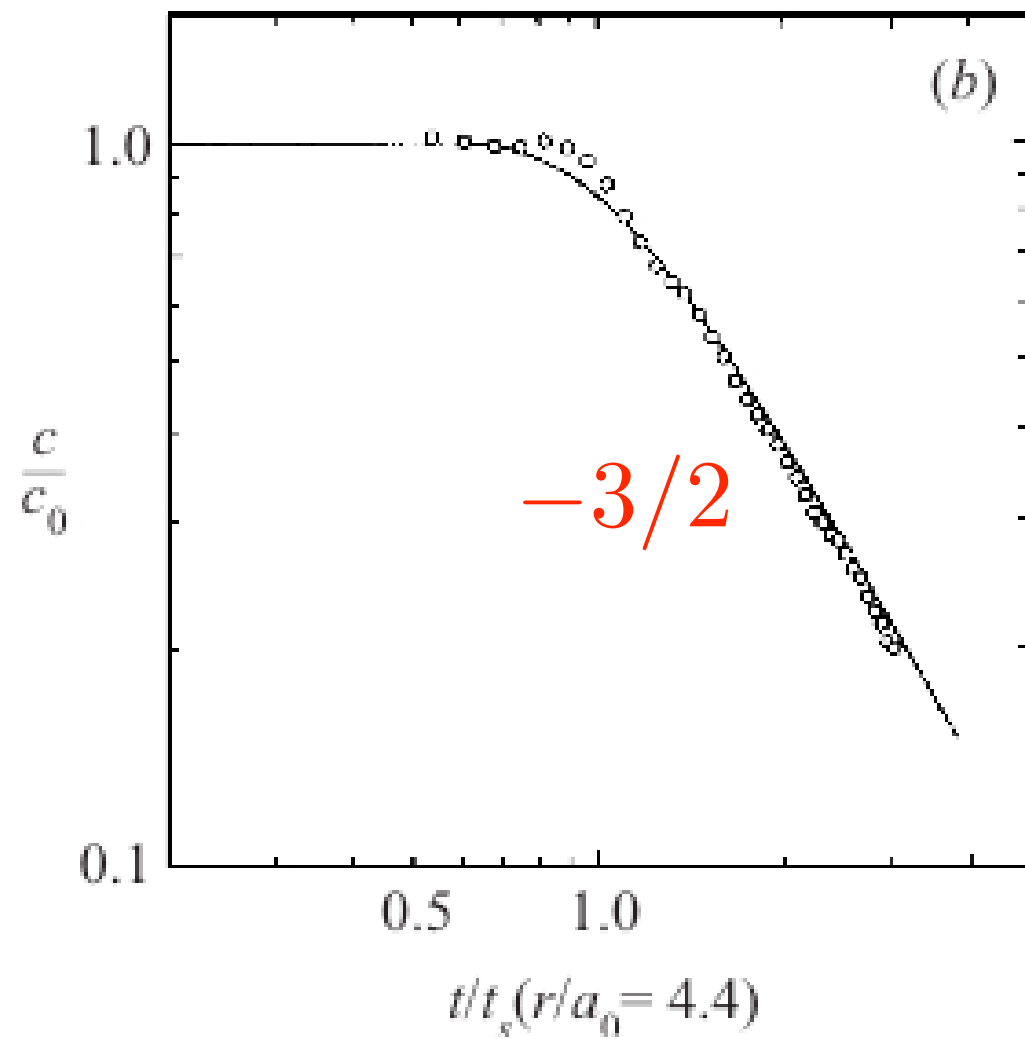
$$C(r, \xi) \approx \operatorname{erf} \left(\frac{1}{4\sqrt{\tau}} \right) e^{-\xi^2/2\tau}$$

With: $\tau(r, t) = \frac{Dt}{s_0^2} + \frac{D\Gamma^2 t^3}{3\pi^2 r^4 s_0^2}$

$$C_{\max}(r, t) \sim t^{-3/2} \quad \text{for } t > t_s$$



fixed times



fixed location

Spatial concentration field

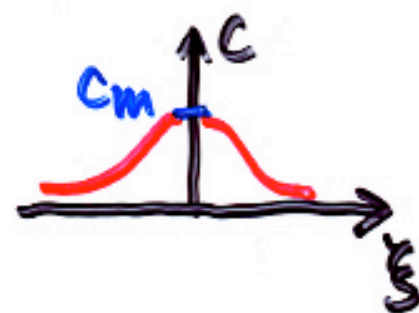
$$c(r, \xi, t) \approx c_0 \operatorname{erf}\left(\frac{1}{4\sqrt{\tau(r)}}\right) \exp\left(\frac{-6\xi^2}{1+24\tau(r)}\right)$$

Concentration distribution

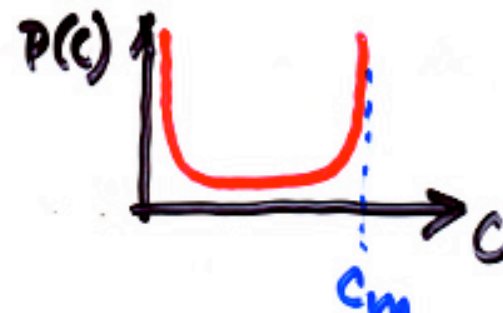
$$P(c) = 2s_0 \int_{c_M(r) > c} |\partial c / \partial \xi|^{-1} dr / A \approx \frac{\sqrt{\tau + 1/24}}{\partial c_M / \partial r}$$

(for details, see Meunier & Villermaux, JFM 476, 2003)

with: $\left| \frac{\partial c}{\partial \xi} \right| = \sqrt{\frac{2}{r(c)}} \cdot c \cdot \sqrt{\ln\left(\frac{c_m(r)}{c}\right)}$

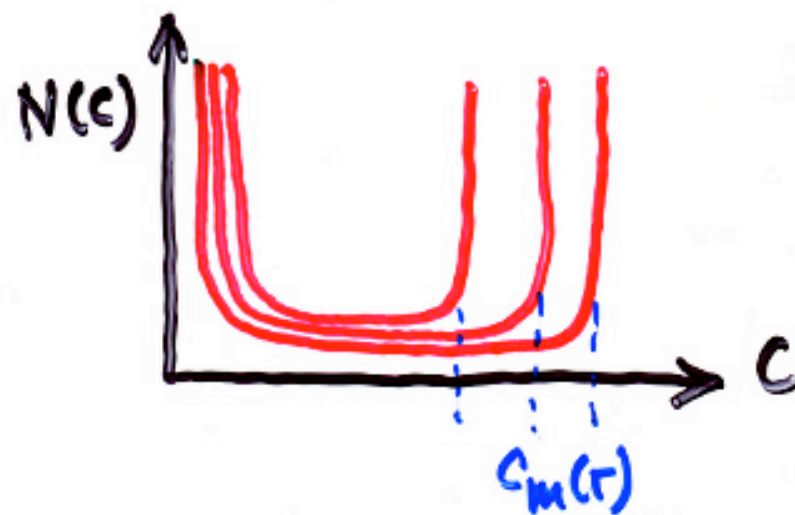


Spatial profile

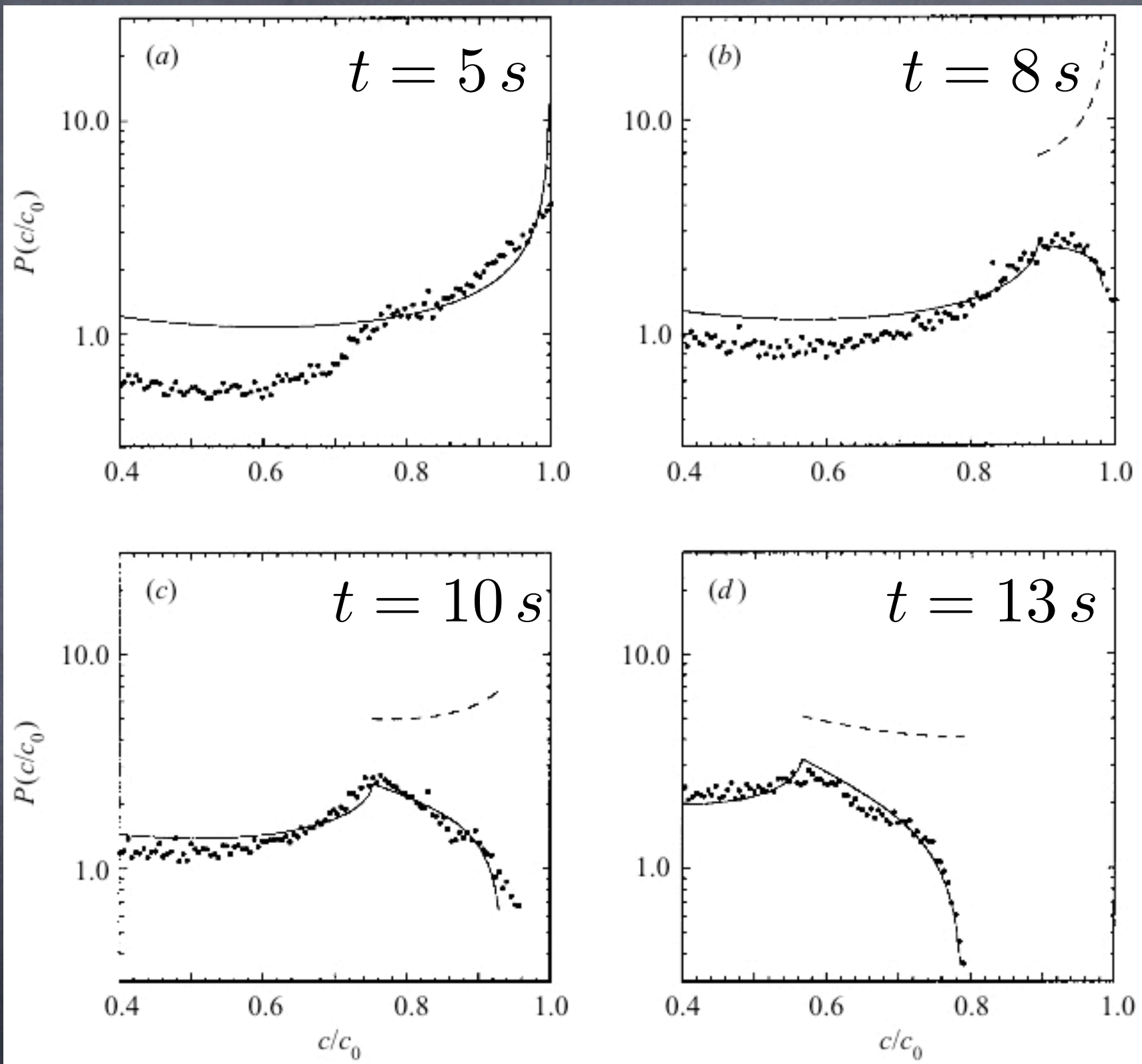


Concentration distribution

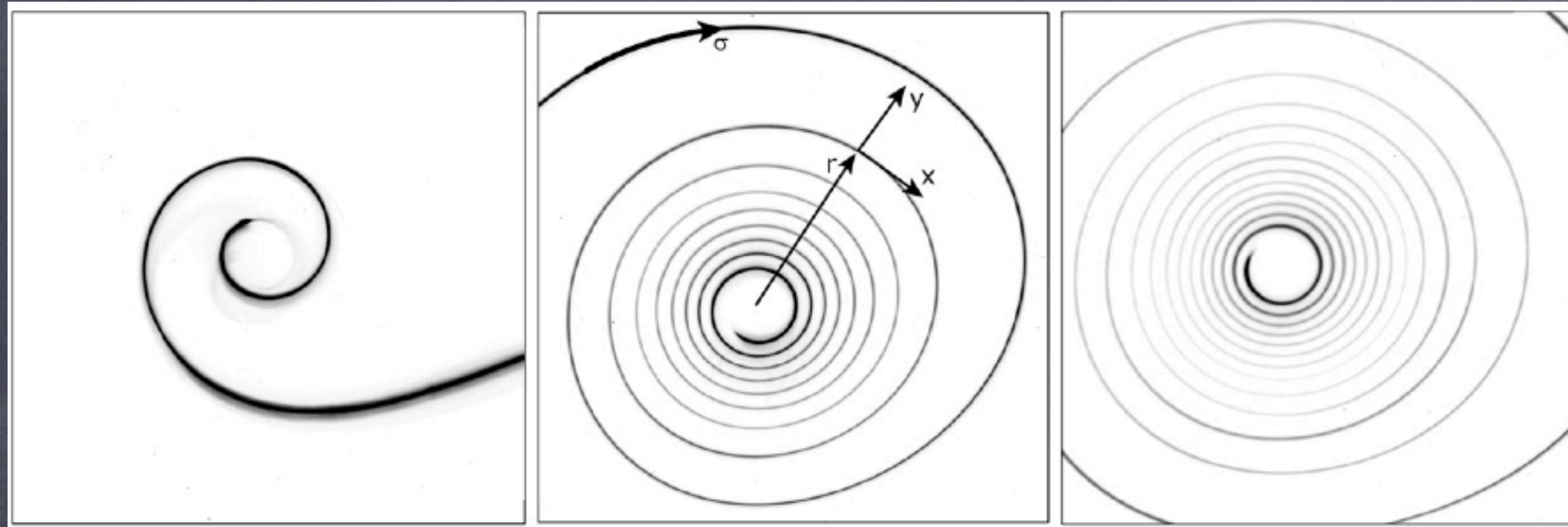
$$N(c)dc = \int d\mu d\eta = \int_{r(c)}^{r_2} \frac{dr dc}{\frac{\partial c}{\partial \xi}}$$



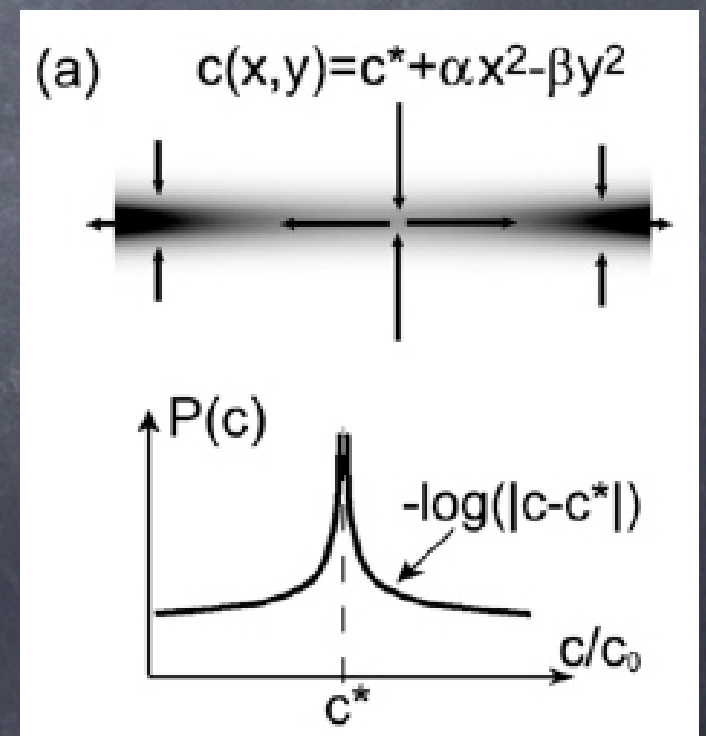
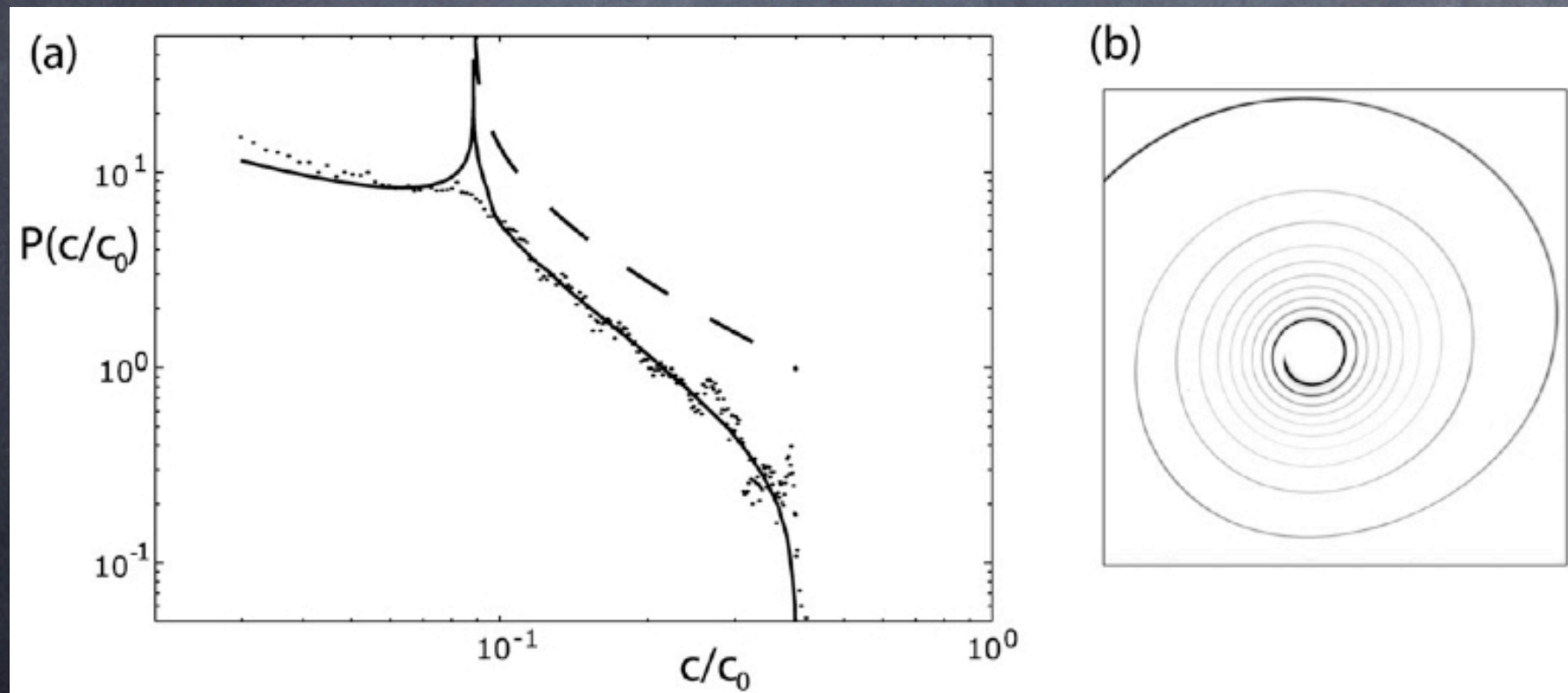
Superposition of U-shaped elementary distributions



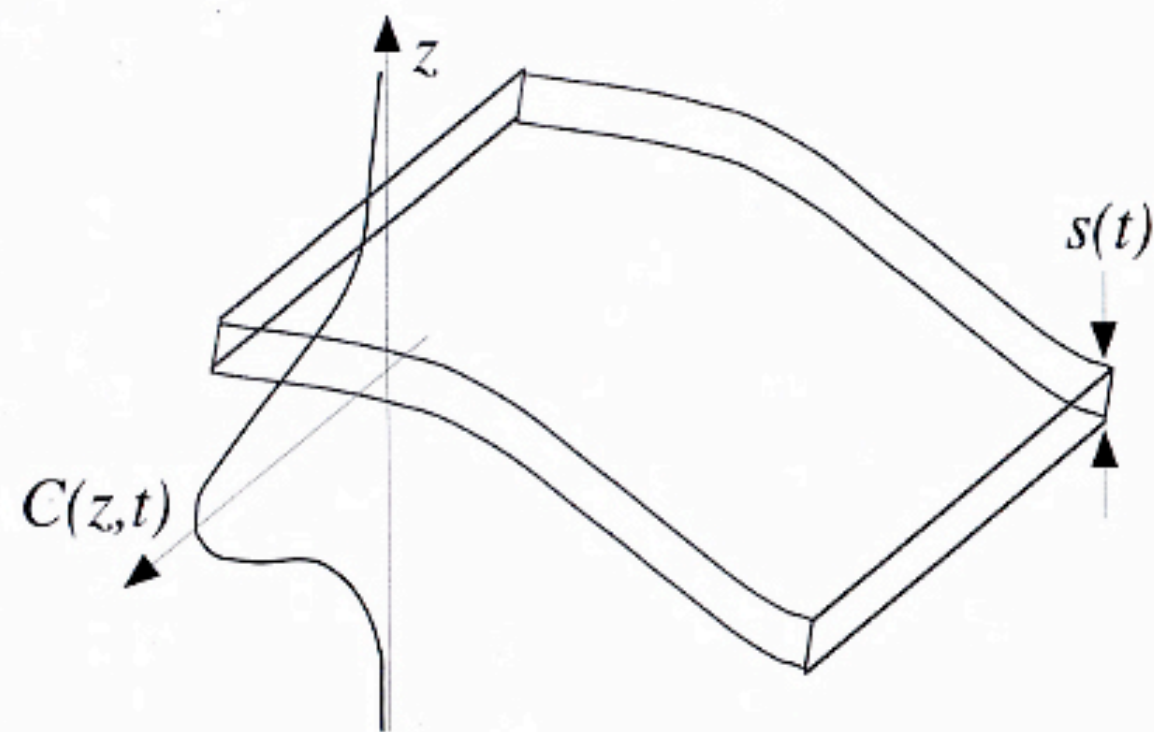
van Hove singularities



spatial saddle
point in $c(x, y)$



Meunier & Villermanx, CRAS 335 (2007)



$$\frac{\partial C}{\partial t} + \frac{ds}{dt} \cdot \frac{\partial C}{\partial z} = D \frac{\partial^2 C}{\partial z^2}$$

Mixing time: $\frac{D}{s^2} \sim \frac{1}{s} \frac{ds}{dt}$

Let: $s = L(\gamma t)^{-\alpha}$, then $\gamma t_s \sim \left(\frac{\gamma L^2}{D}\right)^{\frac{1}{1+2\alpha}}$

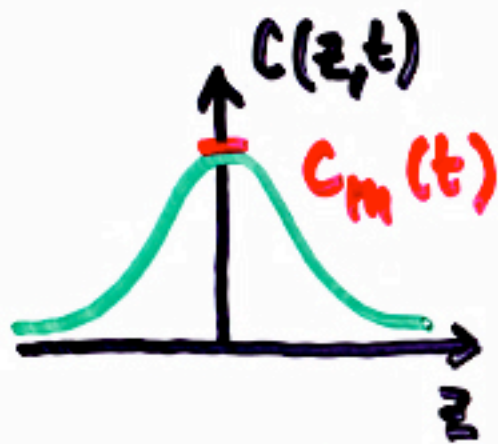
$$Pe = \frac{\gamma L^2}{D} > 1$$

• $\alpha = 1 \rightarrow \frac{1}{1+2\alpha} = \frac{1}{3}$

• $\alpha = 2 \rightarrow \text{"} = \frac{1}{5}$

Let: $s = L e^{-\gamma t}$, then $\gamma t_s \sim \frac{1}{2} \ln\left(\frac{\gamma L^2}{D}\right)$

After the mixing time

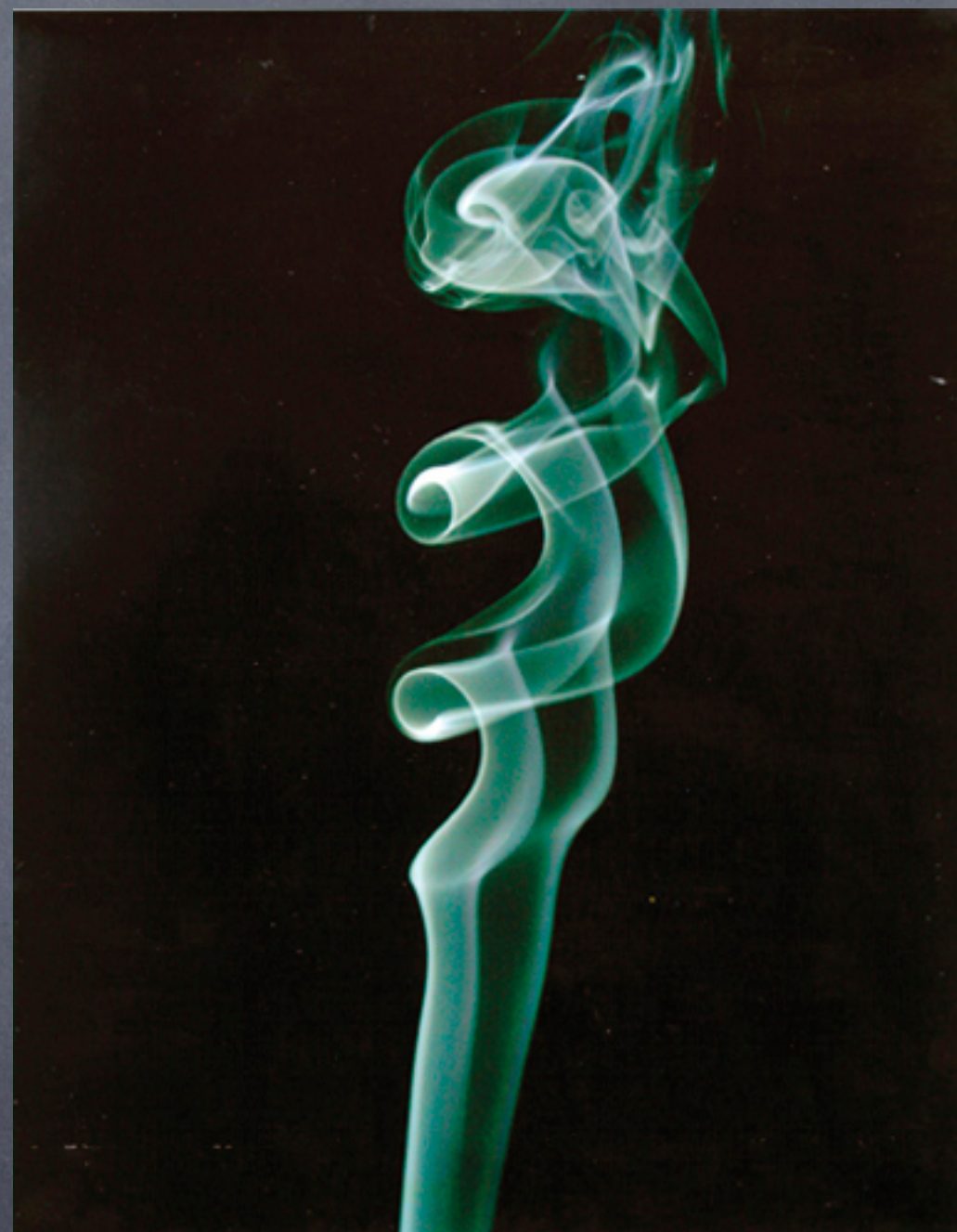


Maximal concentration:

$$C_m(t) \cdot \frac{\sqrt{Dt}}{S(t)} = \text{const.}$$

Thus: $C_m(t) \sim (\gamma t)^{-\frac{1}{2}-\alpha}$

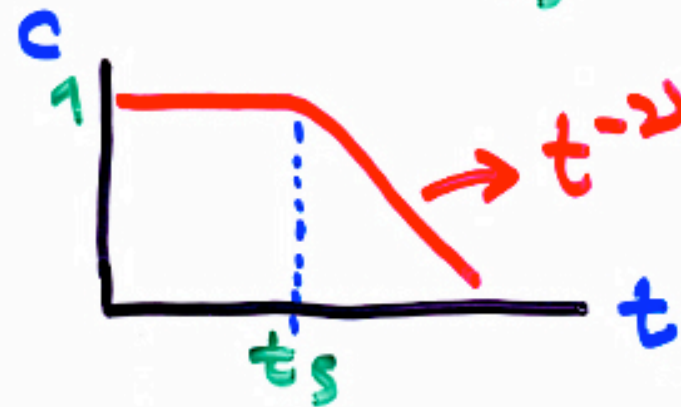
- $\alpha=1 \rightarrow C_m(t) \sim (\gamma t)^{-3/2}$
- $\alpha=2 \rightarrow C_m(t) \sim (\gamma t)^{-5/2}$





- Concentration

$$C = \left(1 + \frac{t}{t_s}\right)^{-\nu}; \quad \nu = \alpha + \frac{1}{2}$$



- Distribution of elongations
 $P(\ell, t)$

Liouville: $P(\ell + \delta\ell, t + \delta t) = P(\ell, t)$

$$\partial_t P = -\partial_\ell \left\{ \frac{\delta\ell}{\delta t} \cdot P \right\}$$

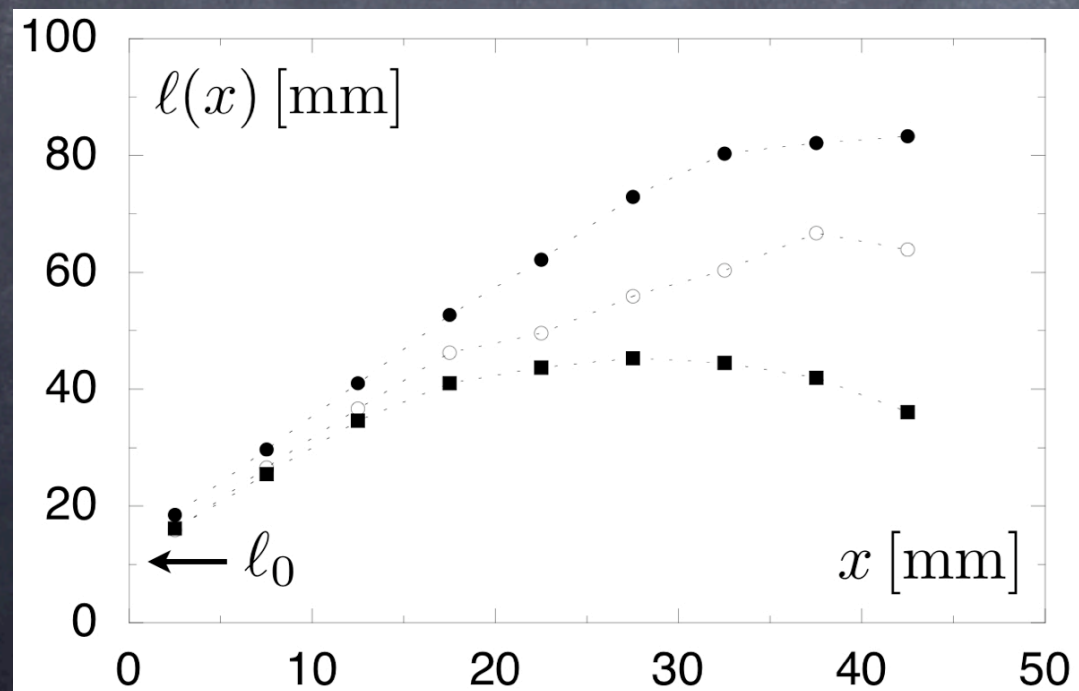
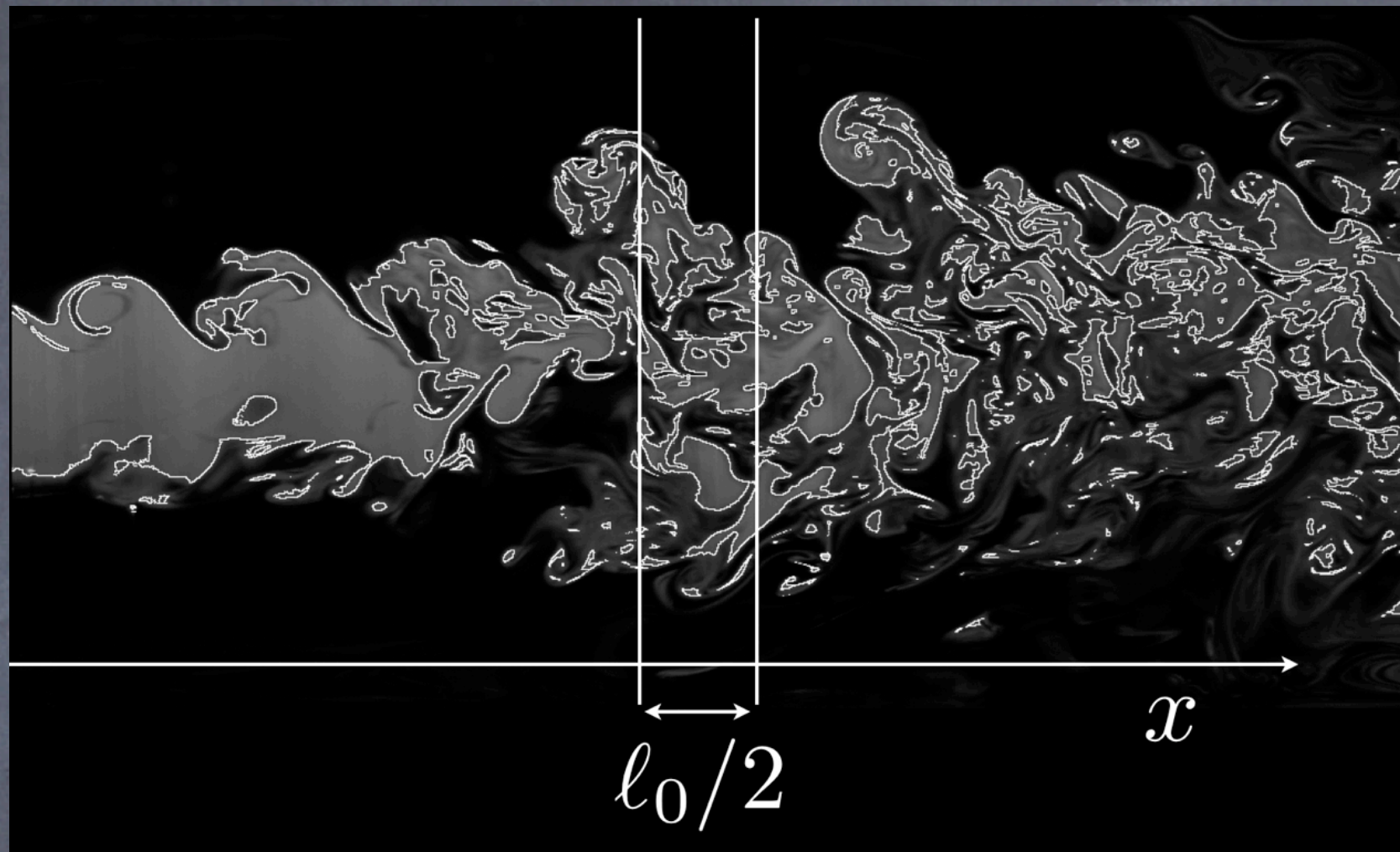
if: $\frac{\delta\ell}{\delta t} = \frac{\ell}{t}$ (linear flows)

then: $\partial_t P = -\partial_\ell \left\{ \frac{\ell}{t} P \right\}$

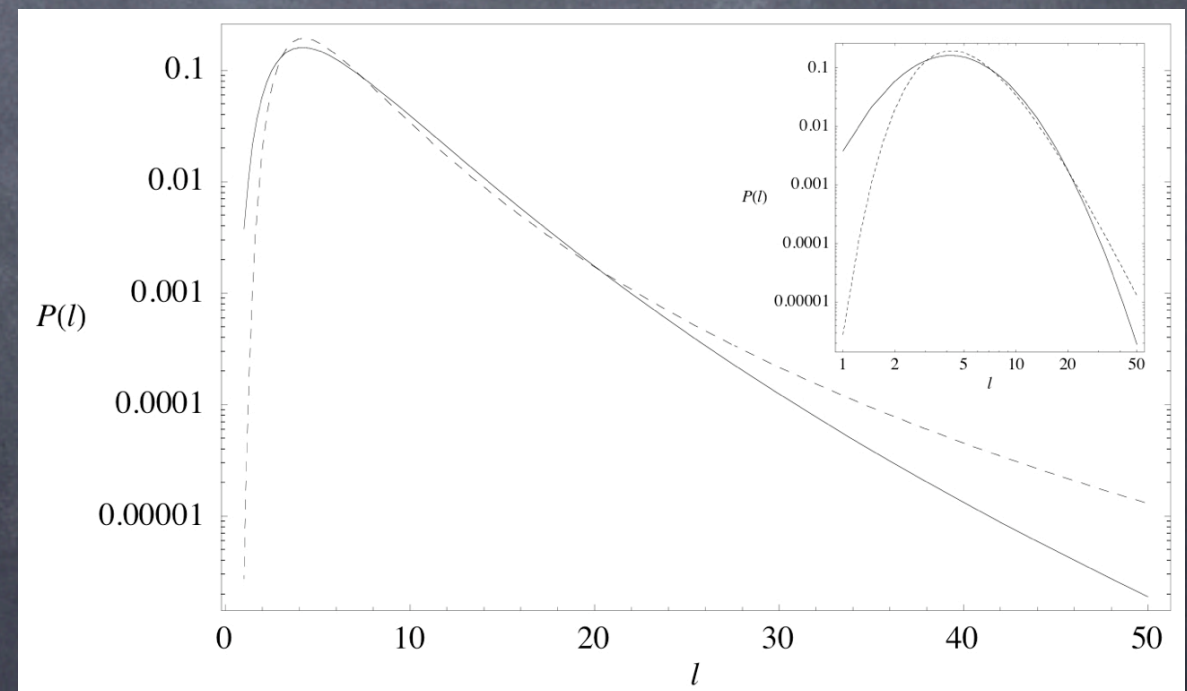
with: $P(\ell, t=0) = \delta(\ell=0)$

$\rightarrow P(\ell, t) = \frac{ut}{\ell^2} \exp\left\{-\frac{ut}{\ell}\right\}$

Contours lengths

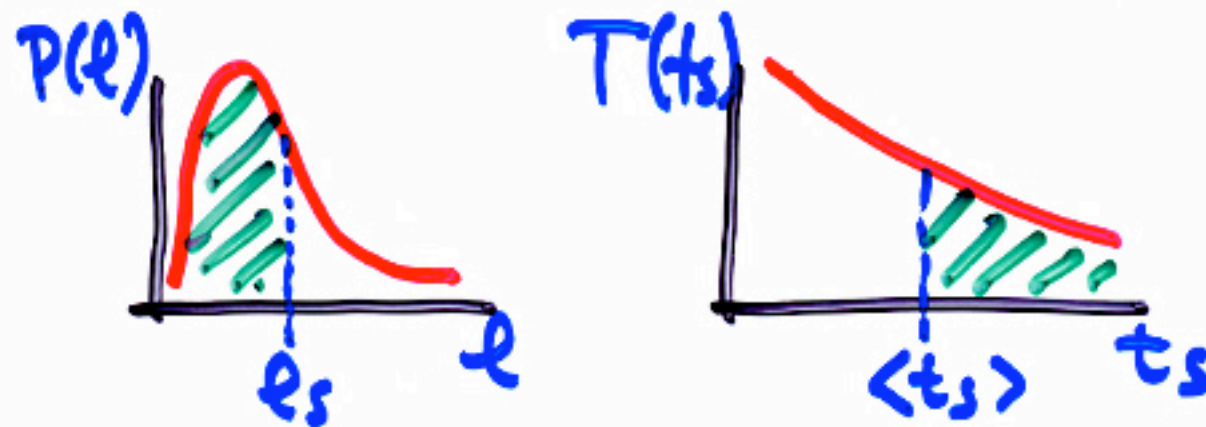


mean



distribution $P(l)$

- Distribution of mixing times



$$\int_{t_s}^{\infty} T(t_s) dt_s = \int_0^{\ell_s} P(\ell, t) d\ell$$

$$T(t_s) = -\partial_{t_s} \int_0^{\ell_s} P(\ell, t) d\ell$$

$$\rightarrow T(t_s) = \frac{1}{\langle t_s \rangle} \exp\left\{-\frac{t_s}{\langle t_s \rangle}\right\}$$

also: Shraiman & Siggia, PRE (1994)

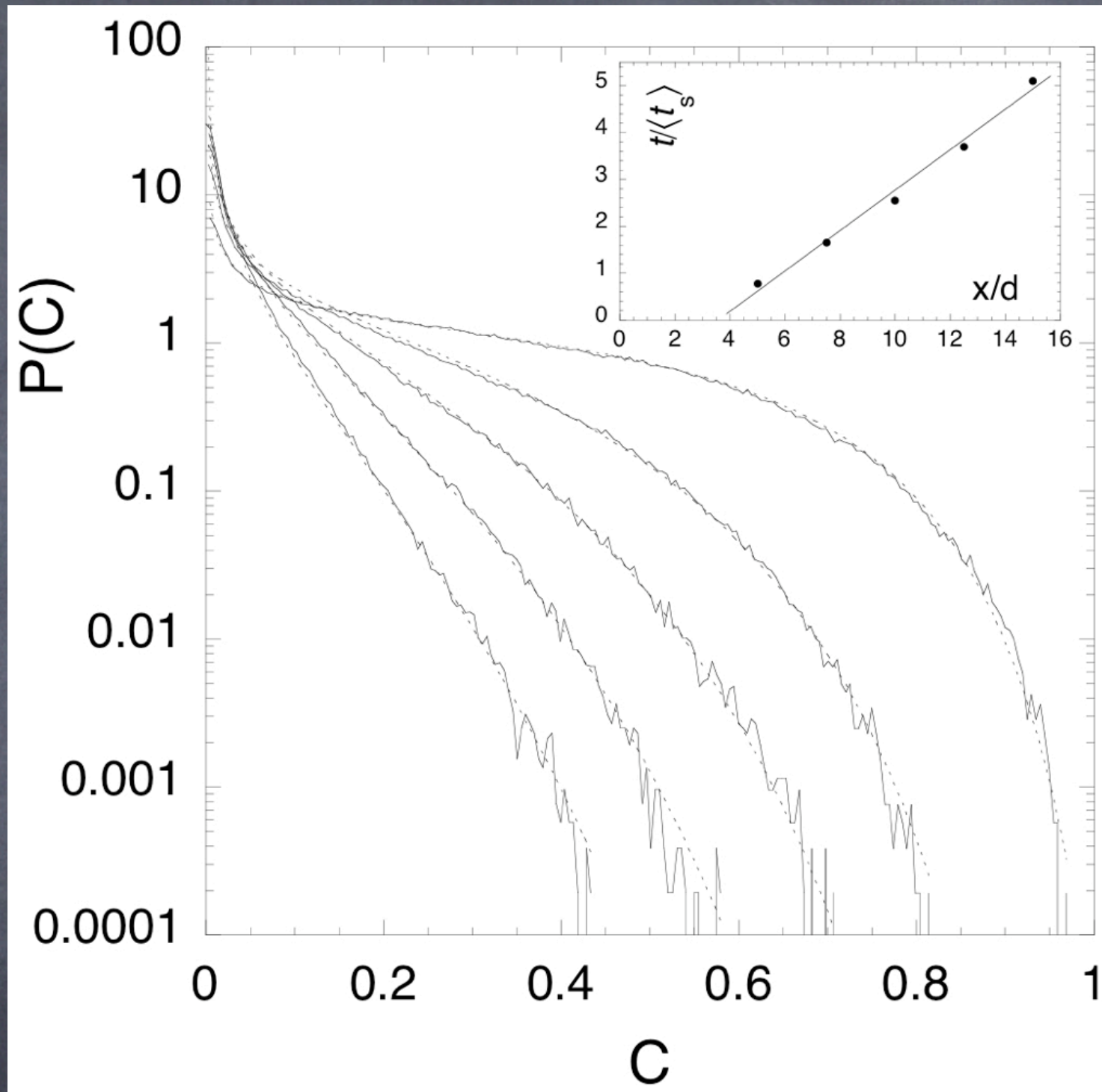
- Distribution of concentration
 $P(c, t)$

$$P(c, t) dc = T(t_s) dt_s$$

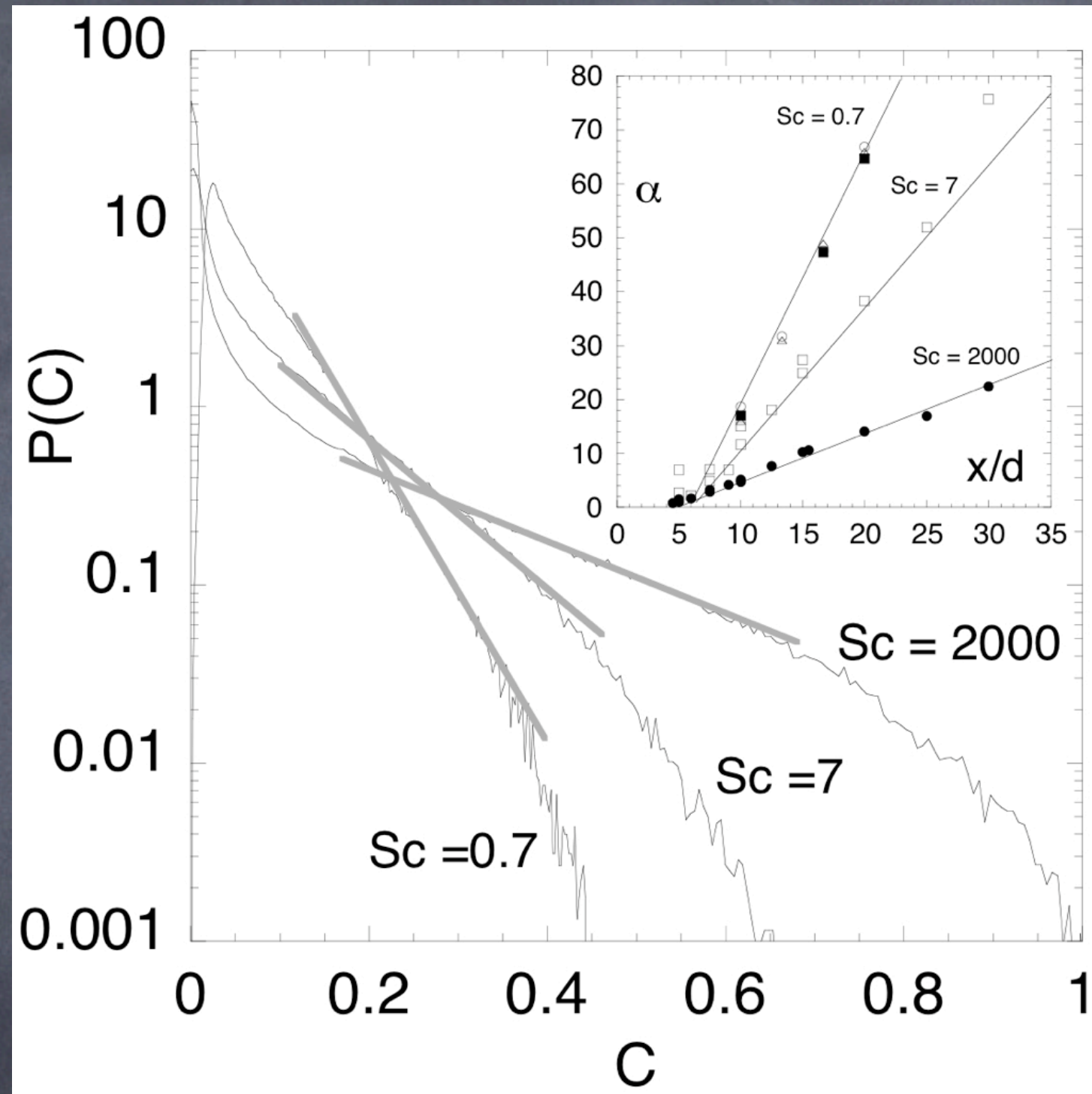
with $c = \left(1 + \frac{t}{t_s}\right)^{-2}$

thus:

$$\left\{ \begin{array}{l} p(c, t) = \frac{\tilde{t}}{2} \frac{c^{-\frac{1+2}{2}}}{(c^{-\frac{1}{2}} - 1)^2} \exp\left\{-\frac{\tilde{t}}{c^{\frac{1}{2}} - 1}\right\} \\ \tilde{t} = \frac{t}{\langle t_s \rangle} \end{array} \right.$$



increasing times, $Sc=7$



fixed time, various scalars



interacting plumes

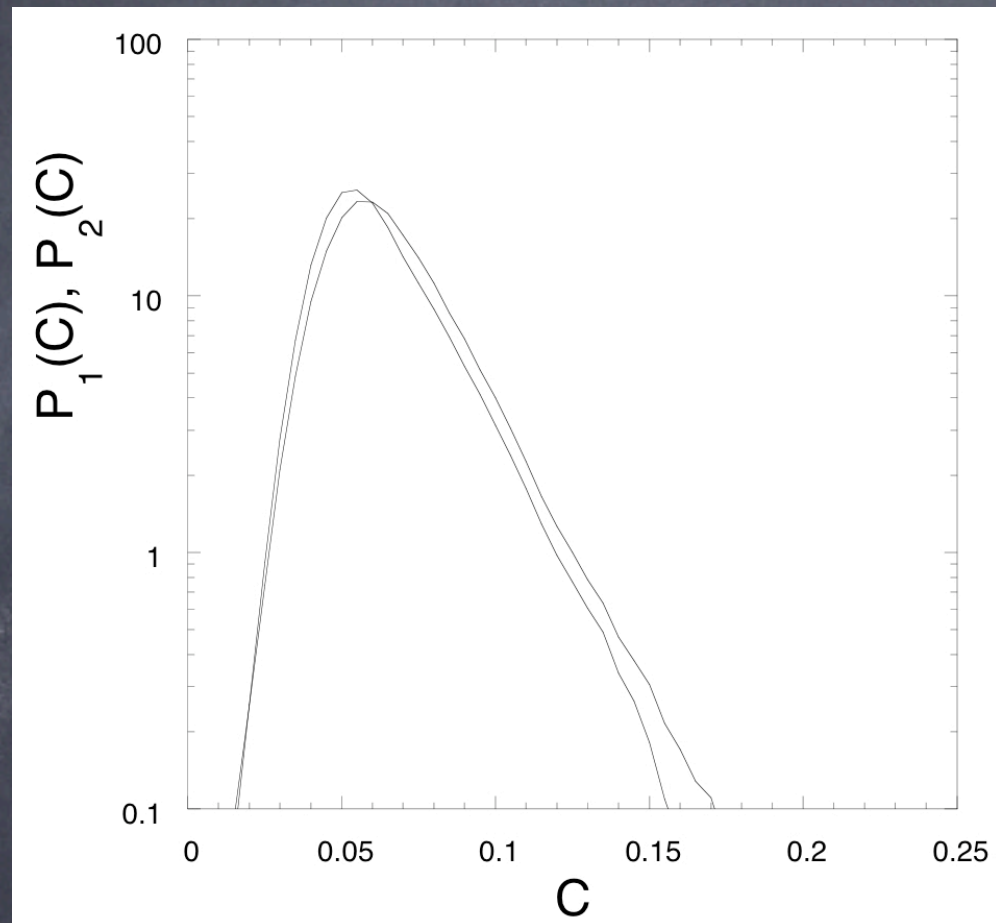
Composition rule

Composition rule

source 1, then source 2

Composition rule

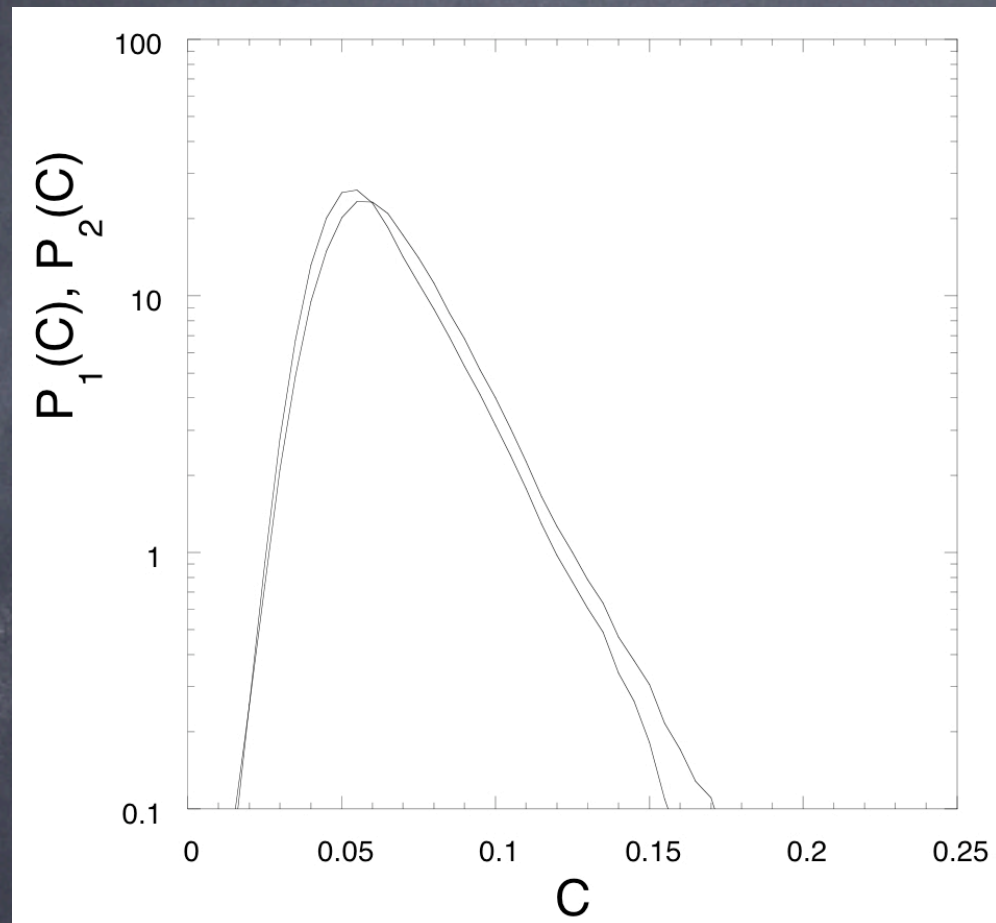
source 1, then source 2



Composition rule

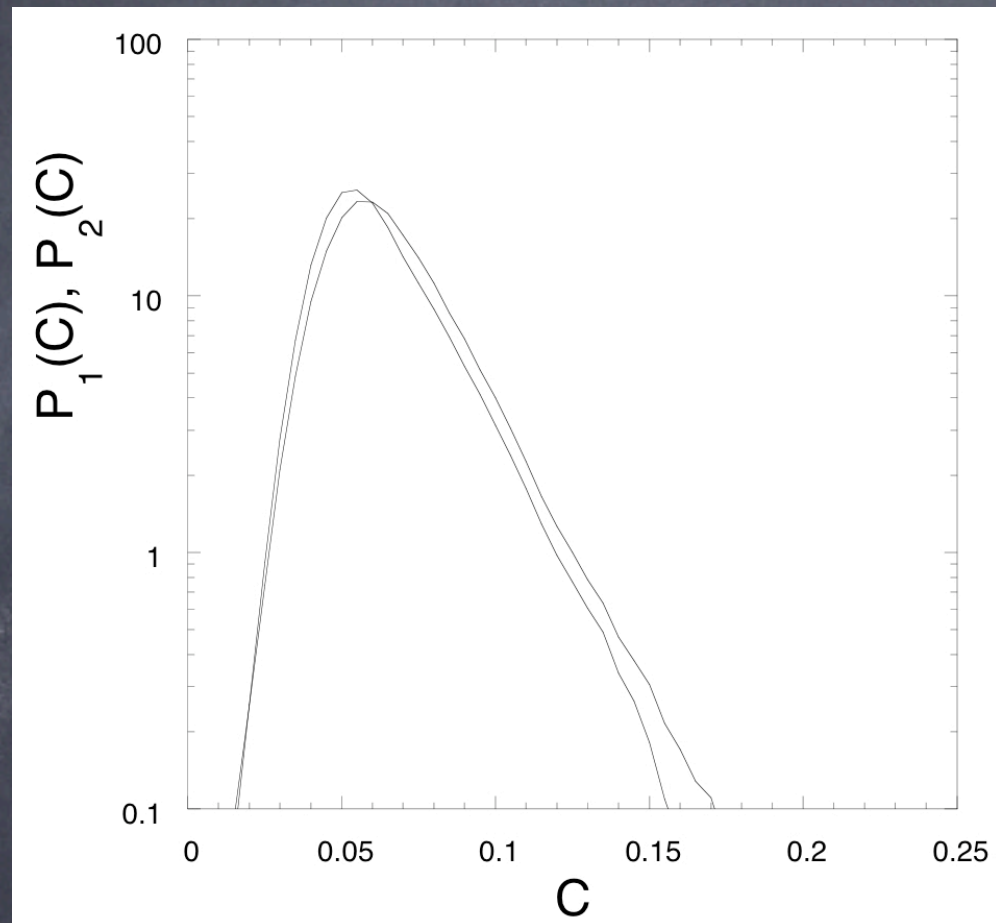
source 1, then source 2

source 1 + source 2

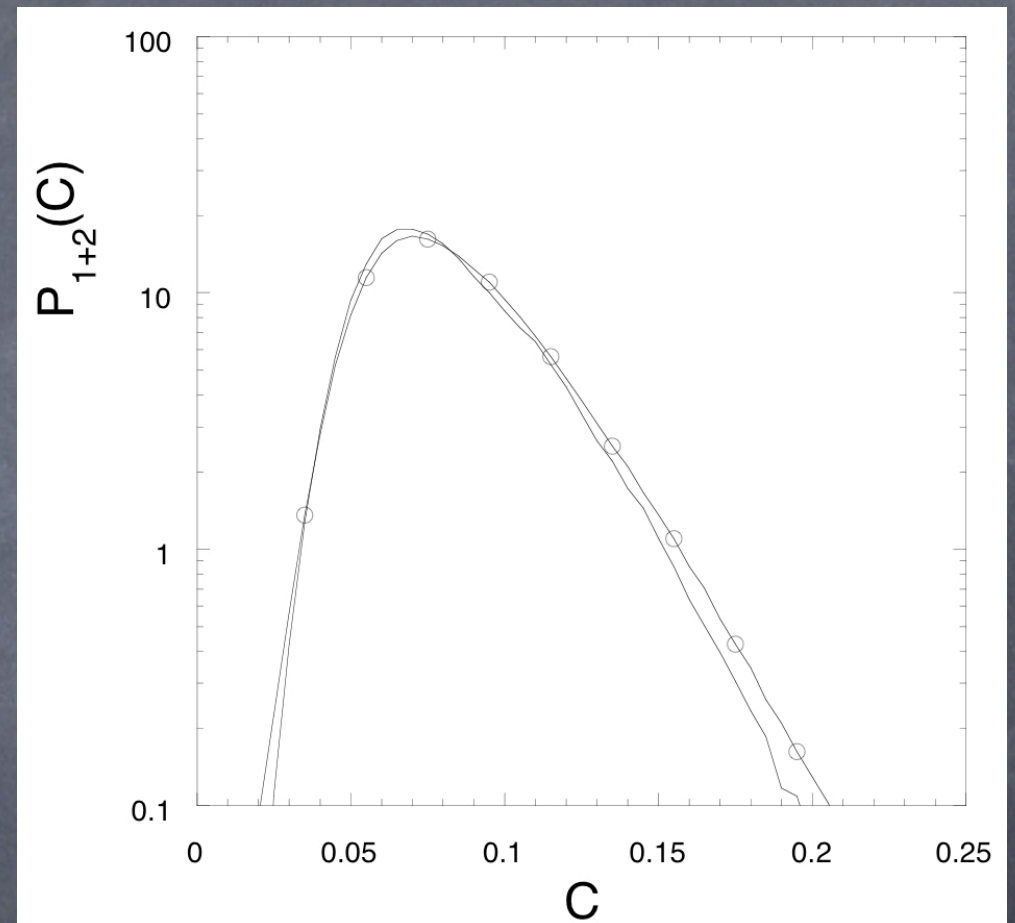


Composition rule

source 1, then source 2

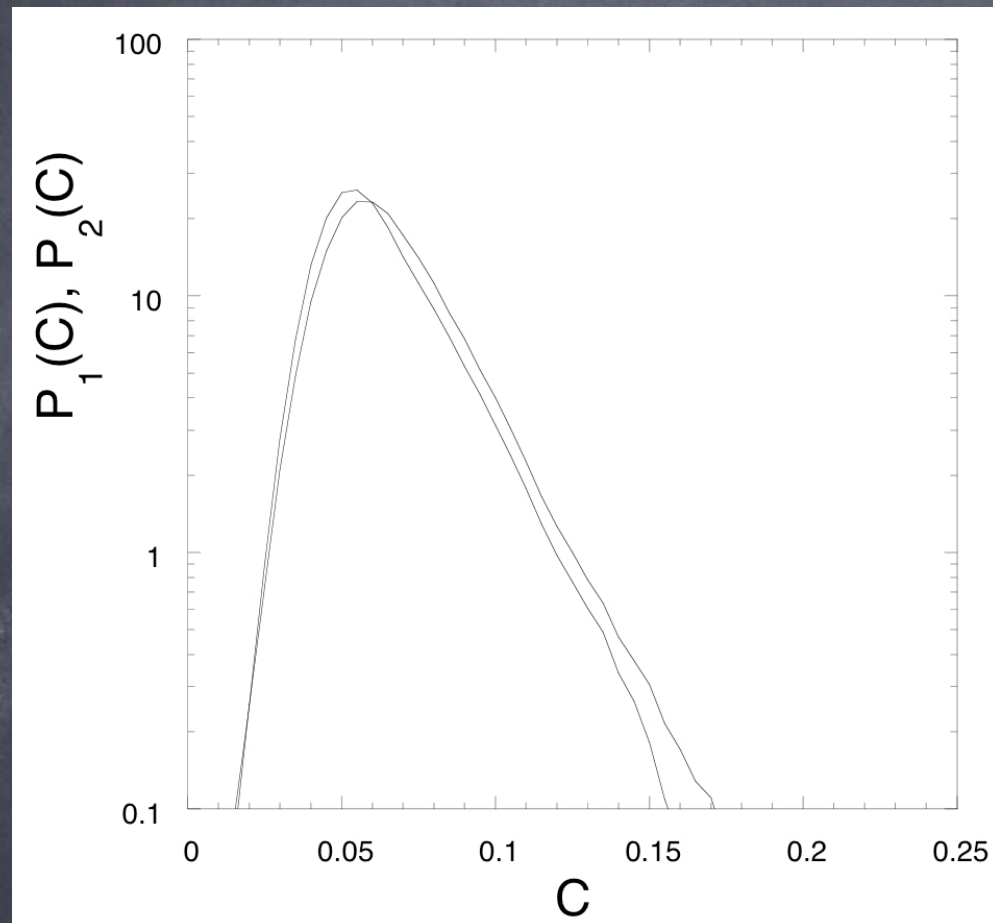


source 1 + source 2

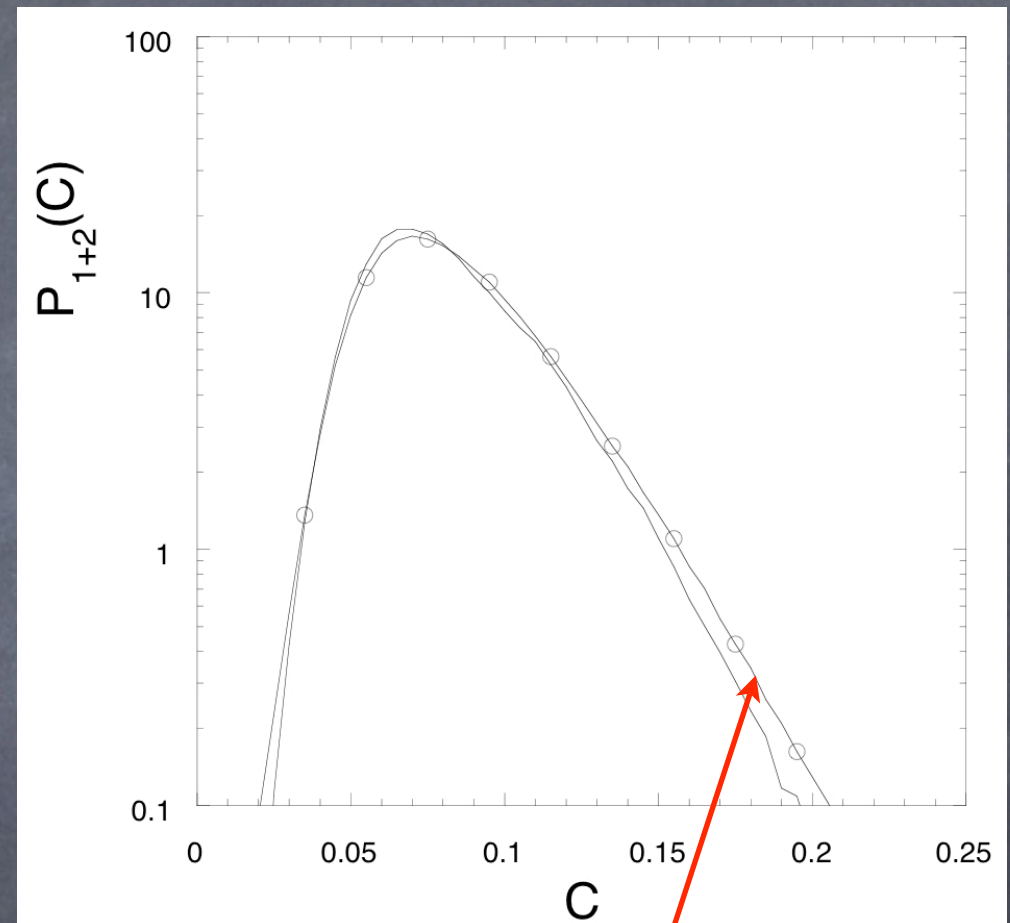


Composition rule

source 1, then source 2



source 1 + source 2



$$P(C) = \int_{C=C_1+C_2} P(C_1)P(C_2)dC_2$$

Microscopic interaction Self convolution processes



$$c = c_1 + c_2$$

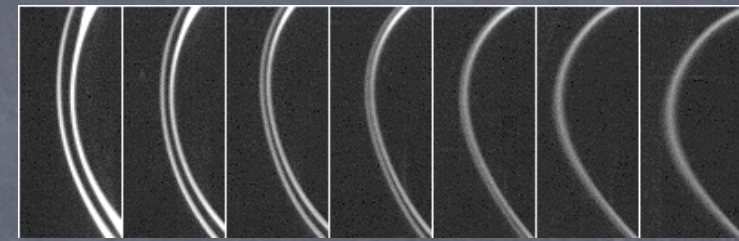
(Fourier)

$$p(c, t + \delta t) = \int_{c=c_1+c_2} p(c_1, t) p(c_2, t) dc_2$$

Convolution

A useful tool: $\tilde{p}(s, t) = \int_0^\infty p(c, t) e^{-s c} dc$
Laplace transform

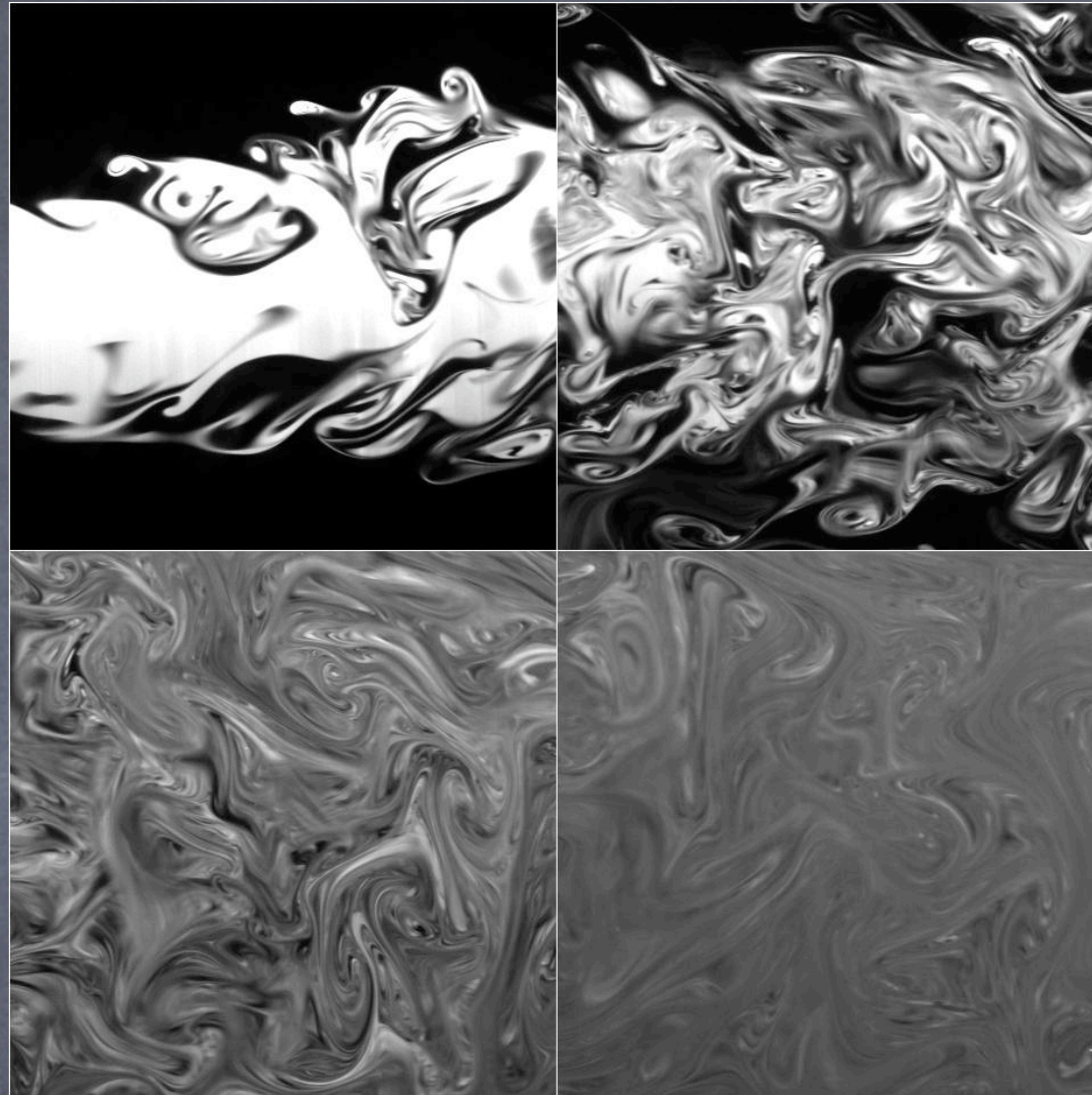
$$p(c, t) \otimes p(c, t) \longleftrightarrow \tilde{p}(s, t)^2$$



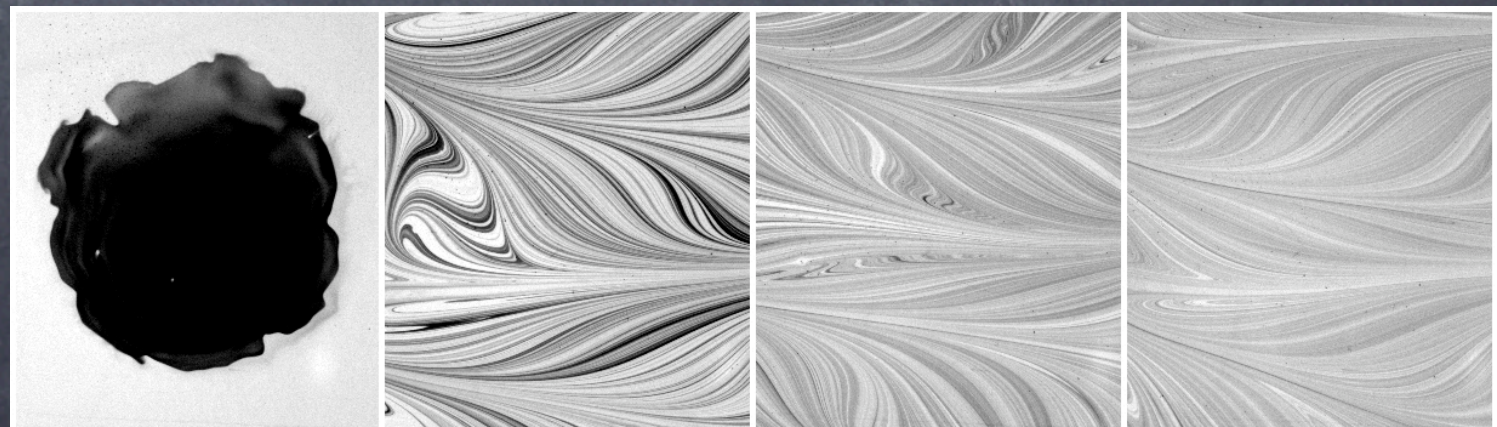
For real !

confined mixtures

turbulent,
self stirred



laminar,
externally stirred



Possible scenarios:

$$(1) \quad \tilde{p}(s, t + \delta t) = \varepsilon \tilde{p}(s, t)^2 + (1 - \varepsilon) \tilde{p}(s, t)$$

$$\text{with } \varepsilon = r \delta t$$

$$\partial_t \tilde{p} = r (-\tilde{p} + \tilde{p}^2)$$

Smoluschowski 1917

$$(2) \quad \tilde{p}(s, t + \delta t) = \tilde{p}^{1 + \varepsilon} \\ = \tilde{p} \exp \{ \varepsilon \ln \tilde{p} \}$$

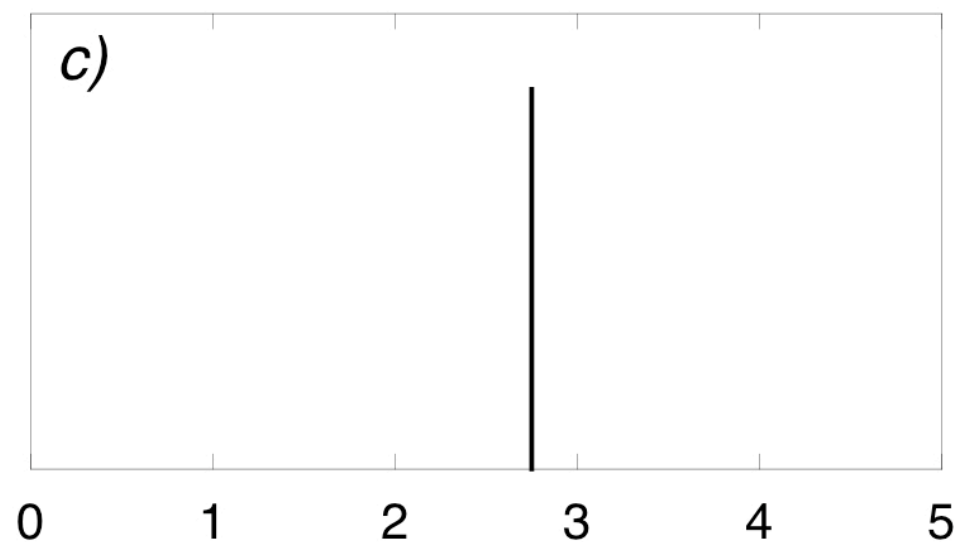
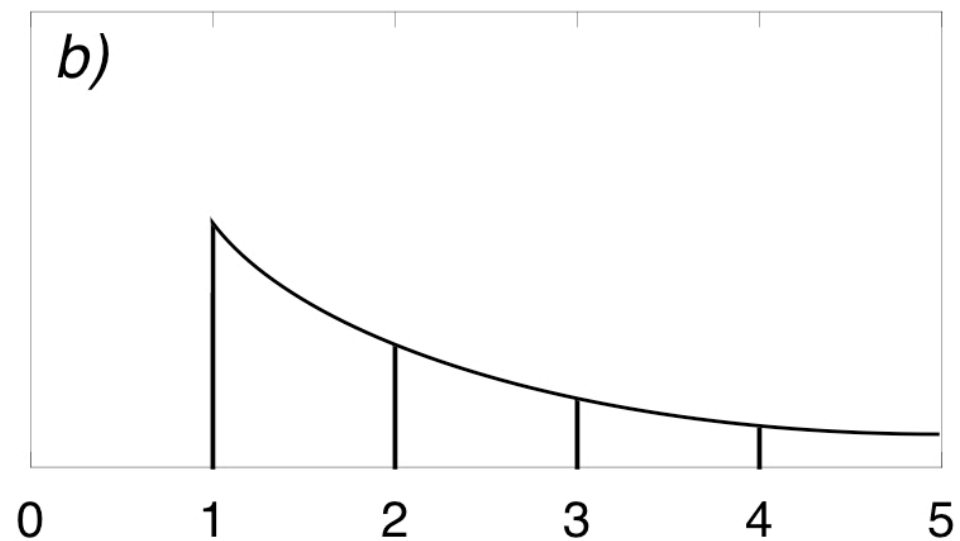
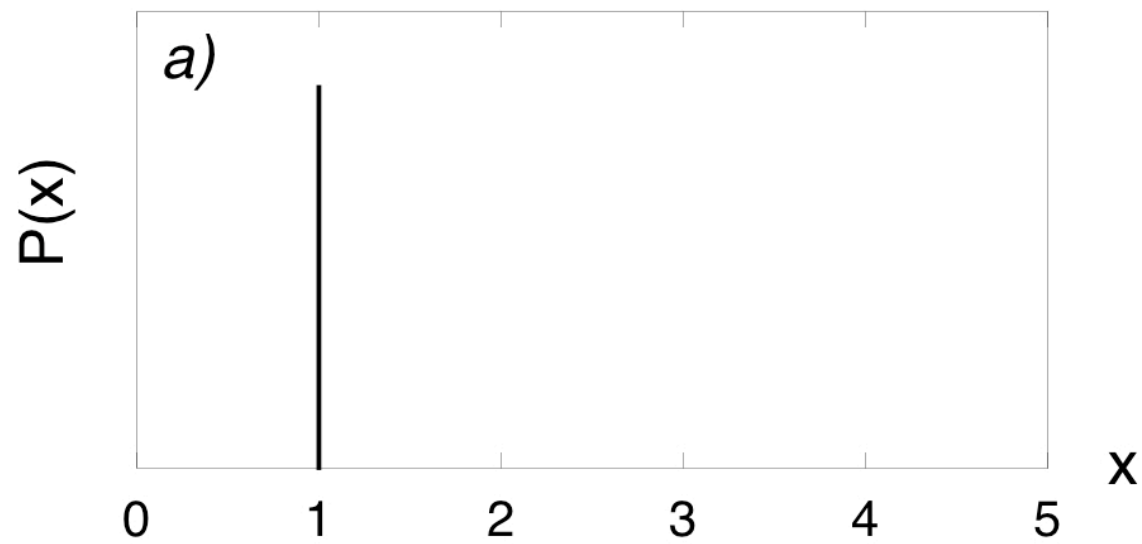
$$\text{with } \varepsilon \rightarrow 0: \partial_t \tilde{p} = r \tilde{p} \ln \tilde{p}$$

(3) General:

$$\partial_t \tilde{p} = r n \left\{ -\tilde{p} + \tilde{p}^{1 + \frac{1}{n}} \right\}$$

$$n = 1 \longrightarrow (1)$$

$$n = \infty \longrightarrow (2)$$



initial condition

$$P(x, t = 0) = \delta(x - 1)$$

$$\partial_t \tilde{P} = -\tilde{P} + \tilde{P}^2$$

$$P(x, t) = e^{-x/e^t} / e^t$$

$$\partial_t \tilde{P} = \tilde{P} \ln \tilde{P}$$

$$P(x, t) = \delta(x - e^t)$$

Plus global shift by stretching:

$$\partial_t p = - \frac{\partial}{\partial c} \left\{ \frac{\delta c}{\delta t} p \right\}$$

$$\frac{\delta c}{\delta t} = -\gamma(t)c ; \gamma(t) = \frac{\alpha + \frac{1}{2}}{t}$$

finally:

$$\partial_t \tilde{p} = -\gamma s \partial_s \tilde{p} + \gamma n \left\{ -\tilde{p} + \tilde{p}^{1+\frac{1}{4}} \right\}$$

Solution:

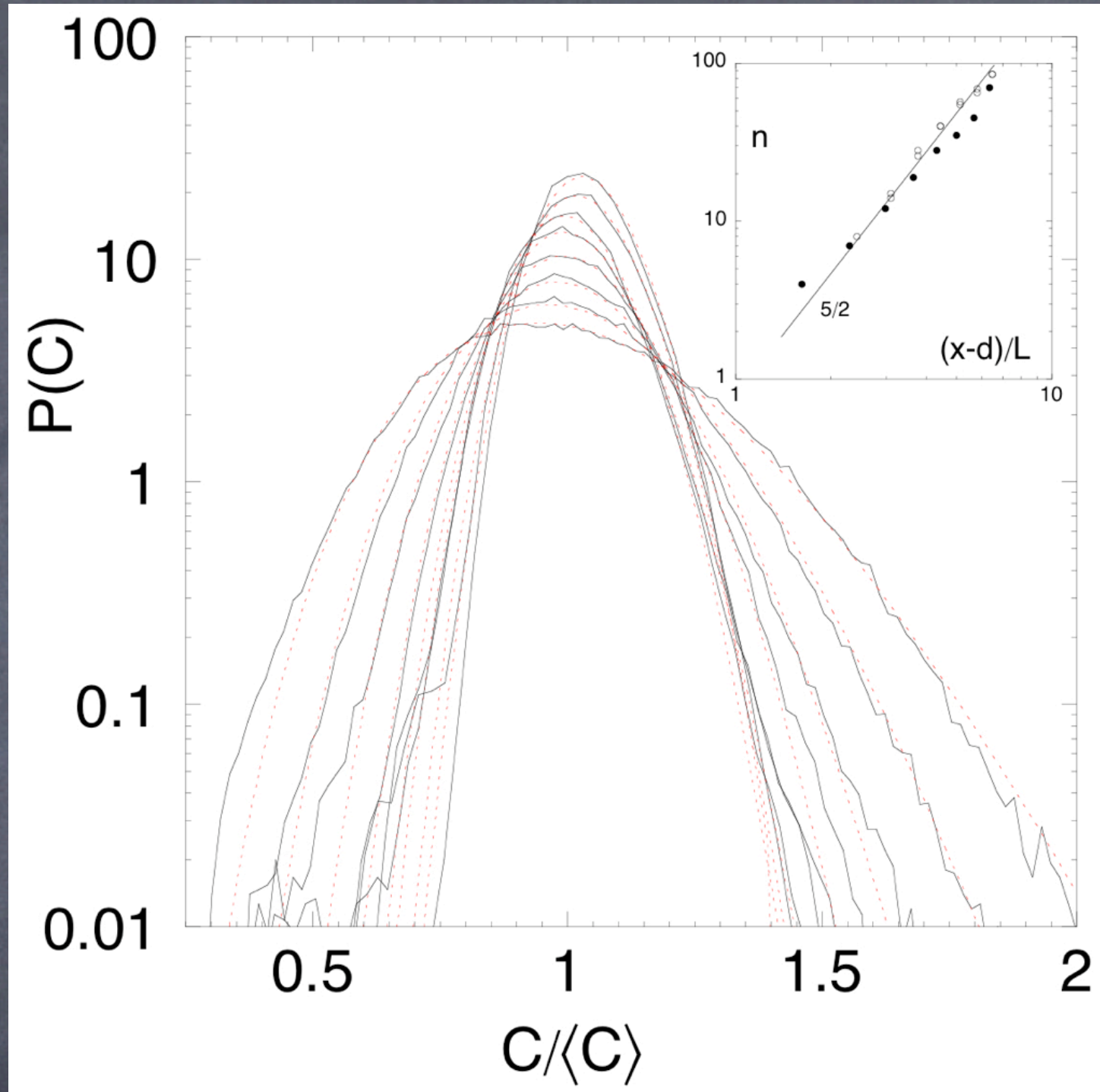
$$\tilde{p} = \left(1 + \langle c \rangle \frac{s}{n} \right)^{-n}$$

$$\text{if } \langle c \rangle = c^{\text{st}} \Rightarrow \gamma = \Gamma ; \frac{dn}{dt} = \gamma n$$

$$\text{Thus: } p(x = \frac{c}{\langle c \rangle}) = \frac{n^n}{\Gamma(n)} x^{n-1} e^{-nx}$$

Gamma Distribution

$$\gamma = \frac{5/2}{t} \longrightarrow n(t) \sim t^{5/2}$$

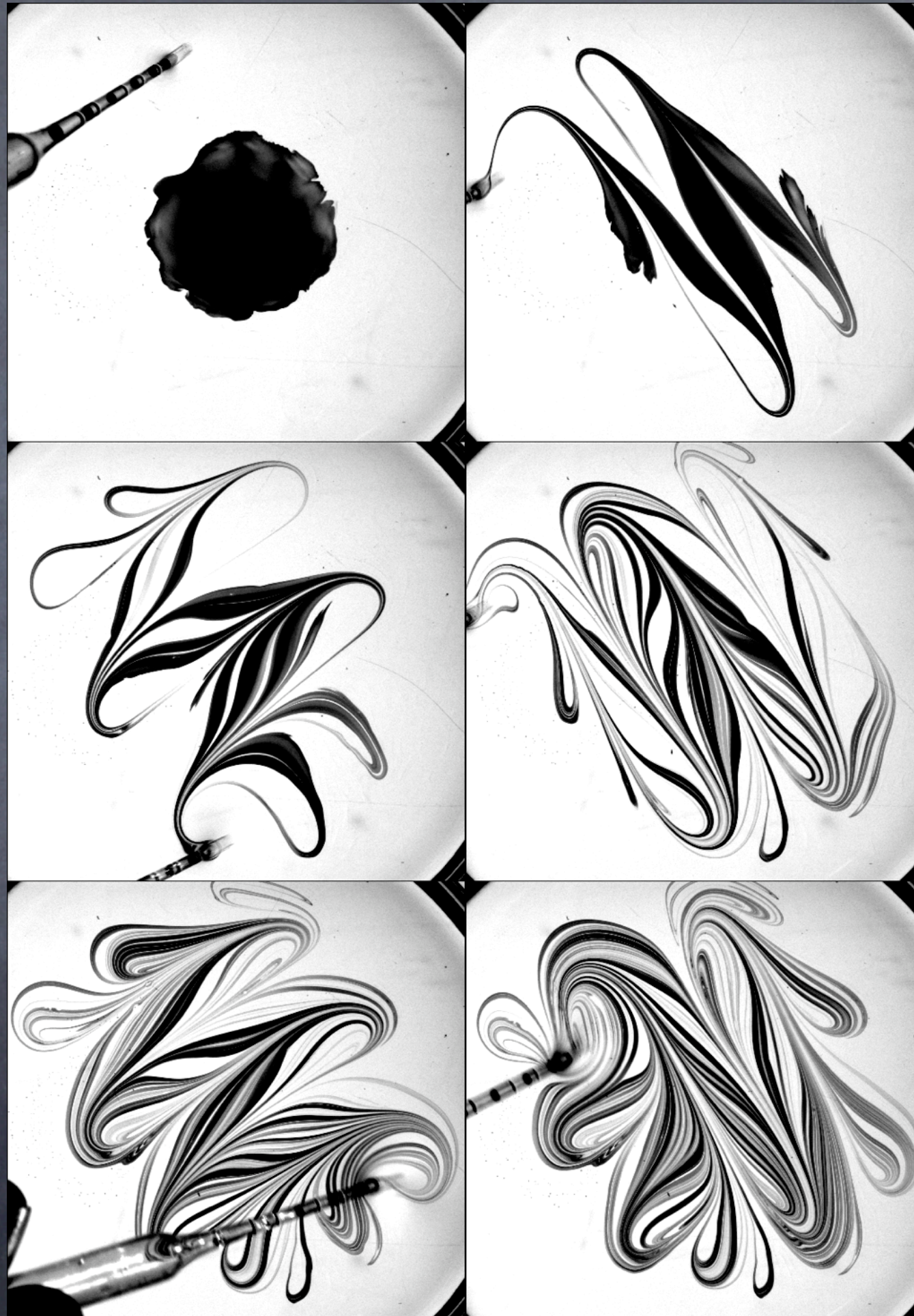


$$n \sim t^{2+1/2}$$

$$P(X = C/\langle C \rangle) = \frac{n^n}{\Gamma(n)} X^{n-1} e^{-nX}$$

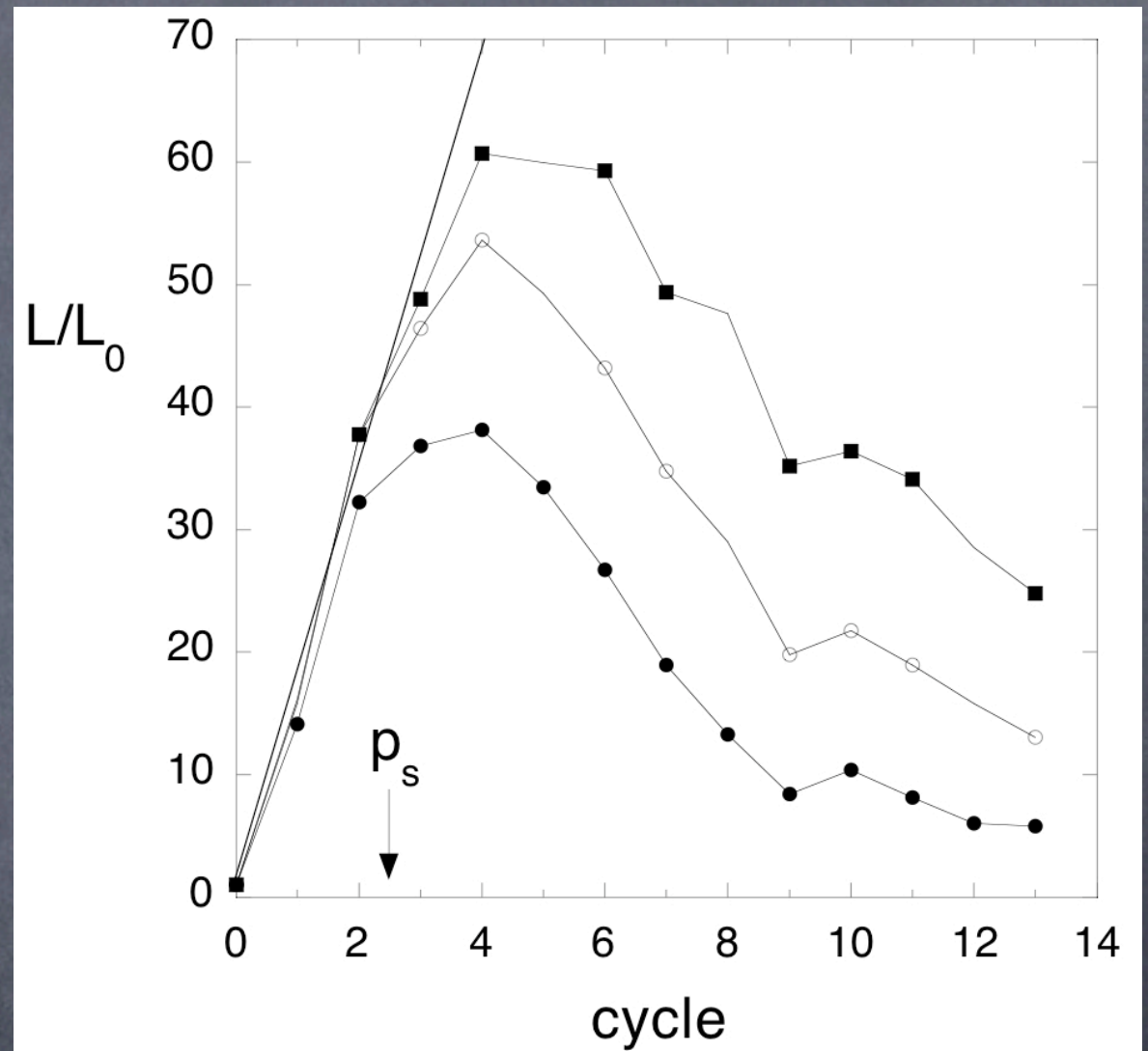
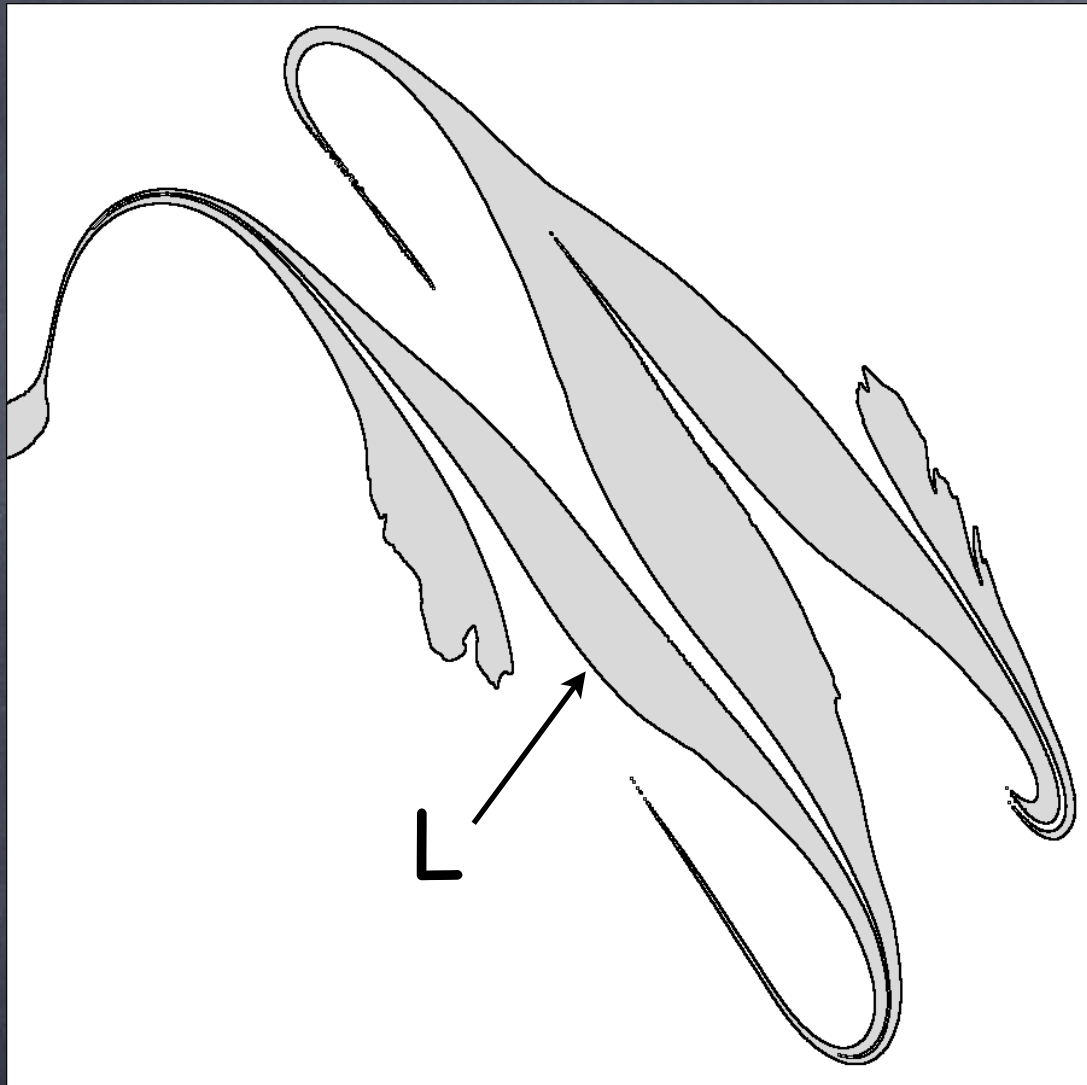
Villermanx & Duplat, PRL 91, 184501 (2003)

Stirring Protocol

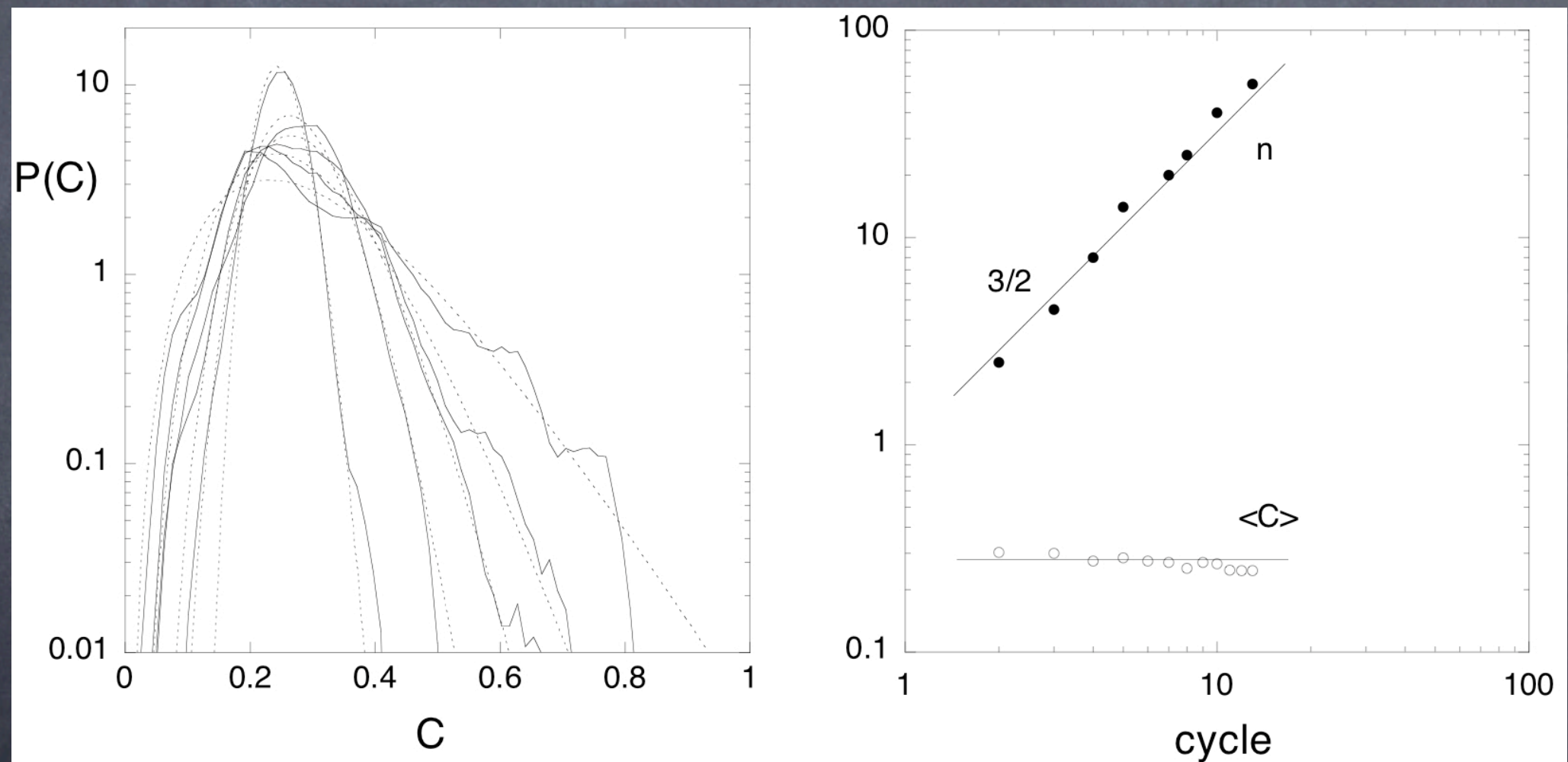
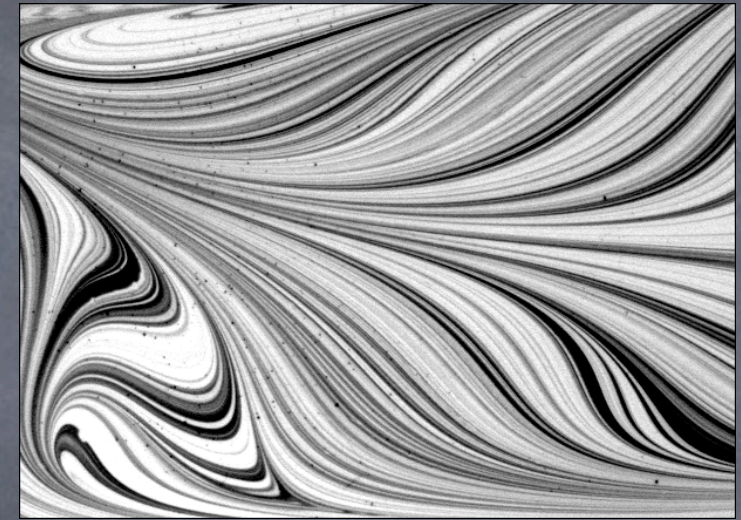


$$Re = 0.1$$

Material Contours lengths



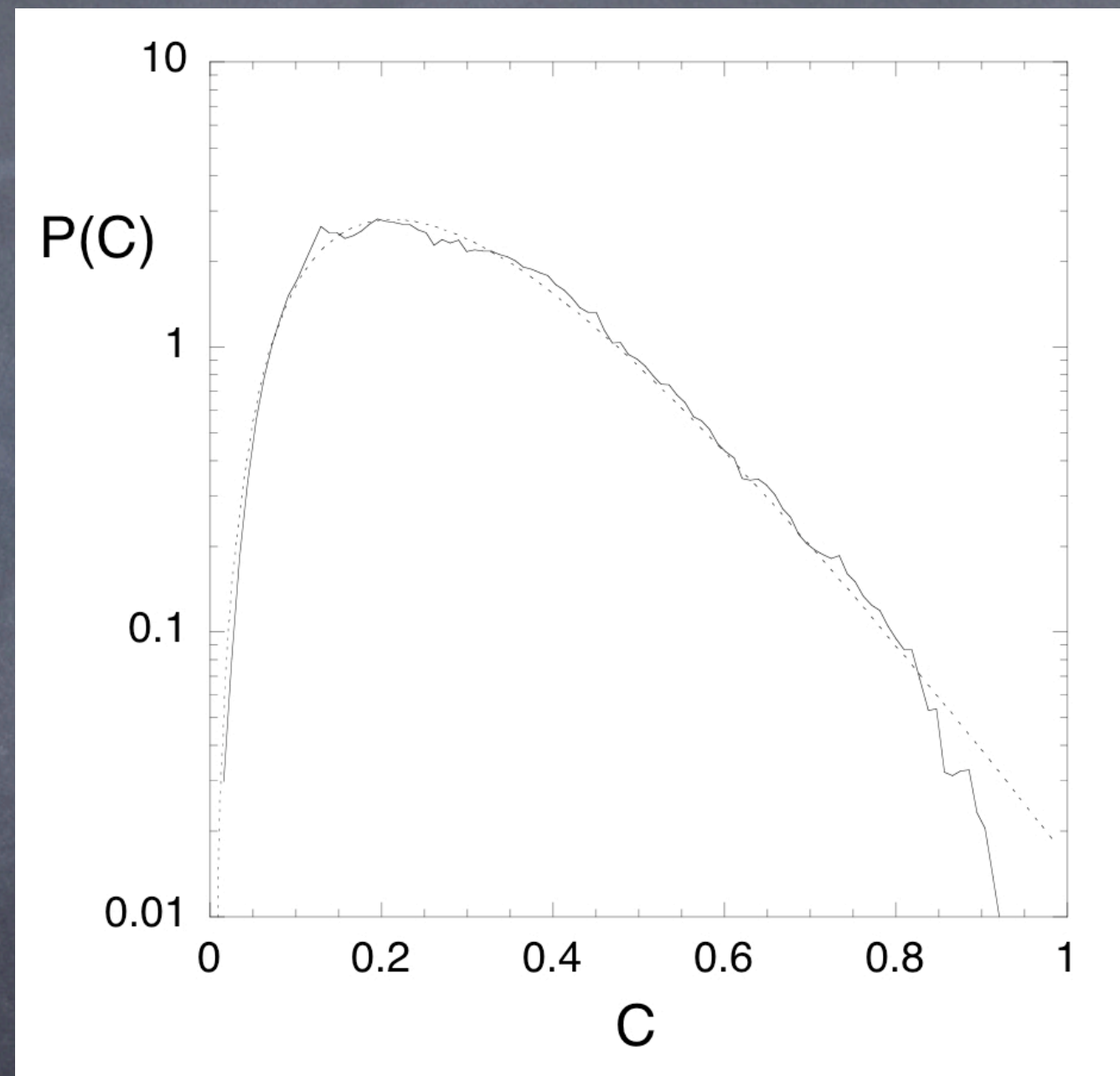
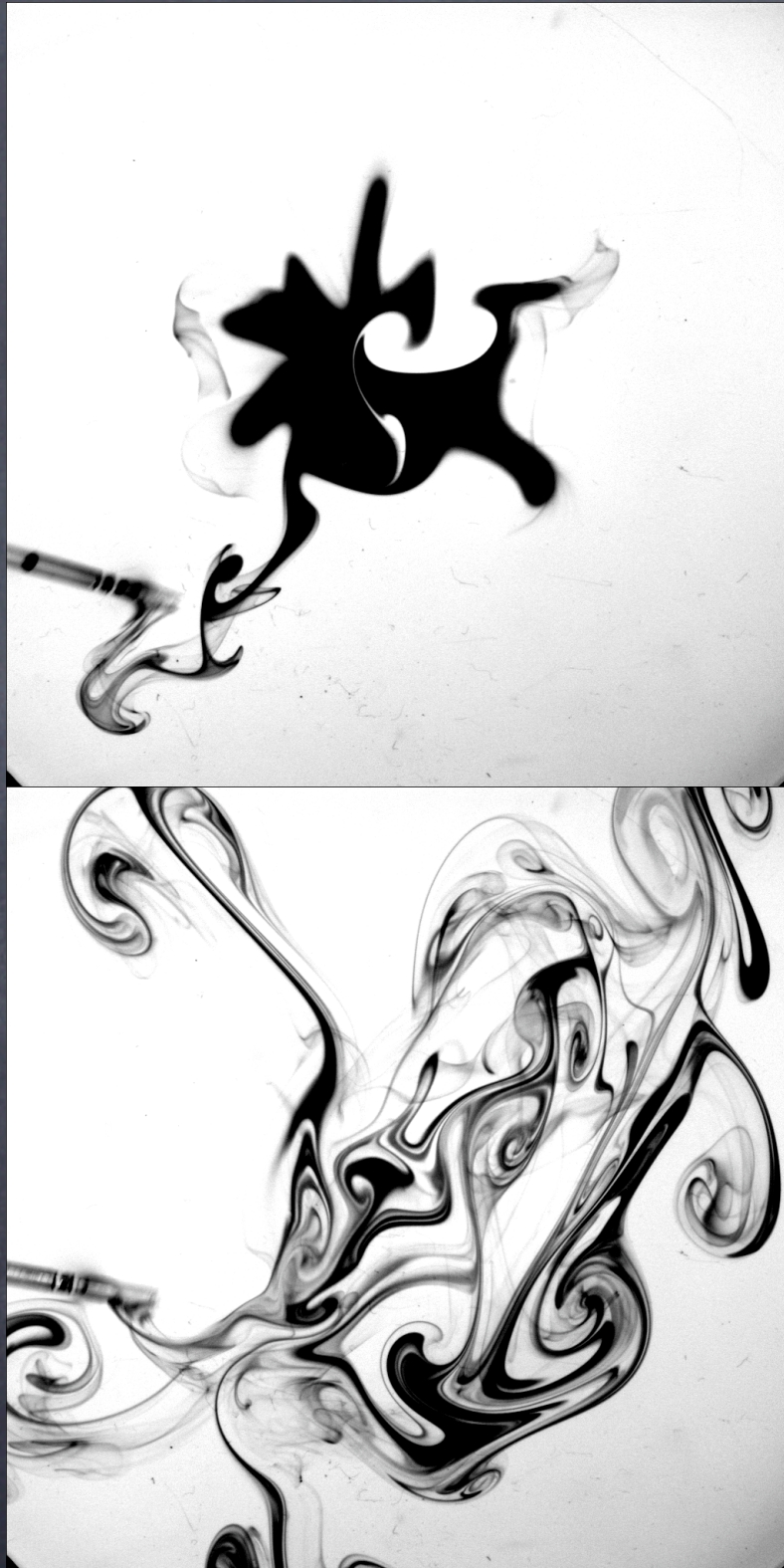
Concentration distribution



$$n \sim (\text{cycle})^{1+1/2}$$

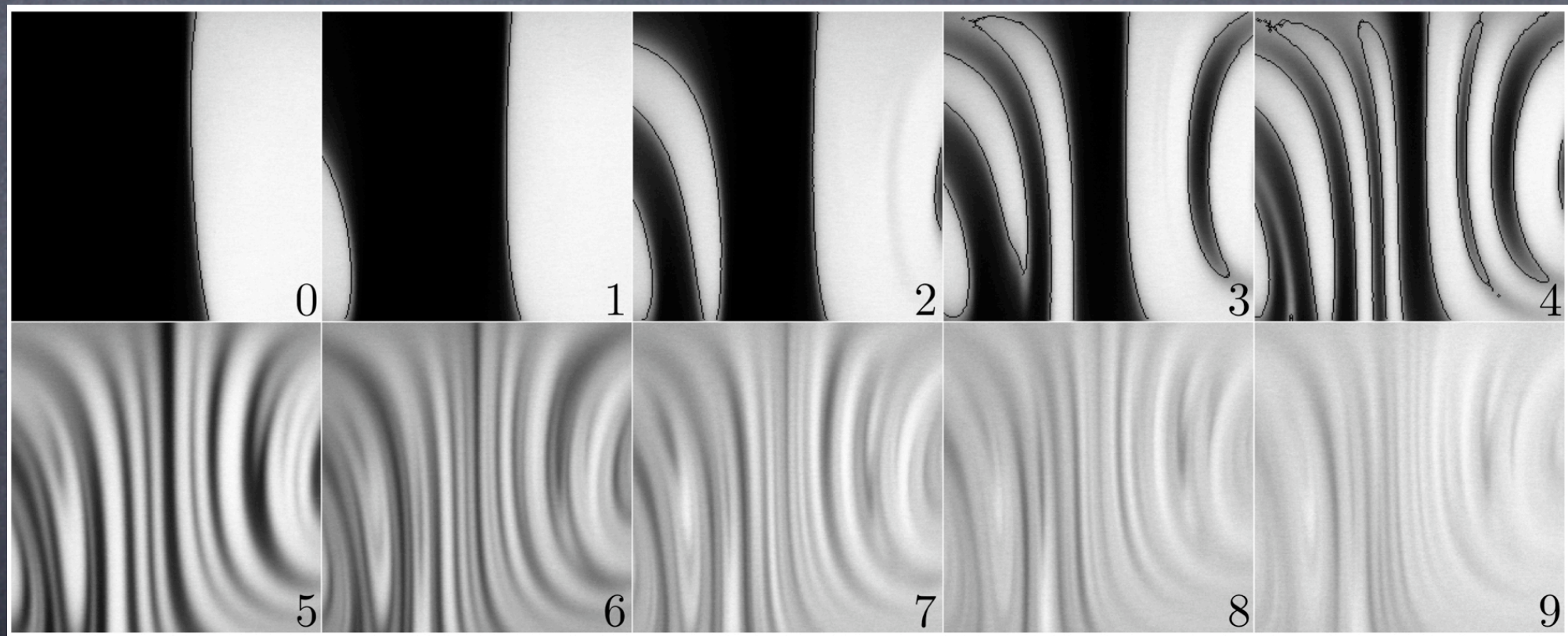
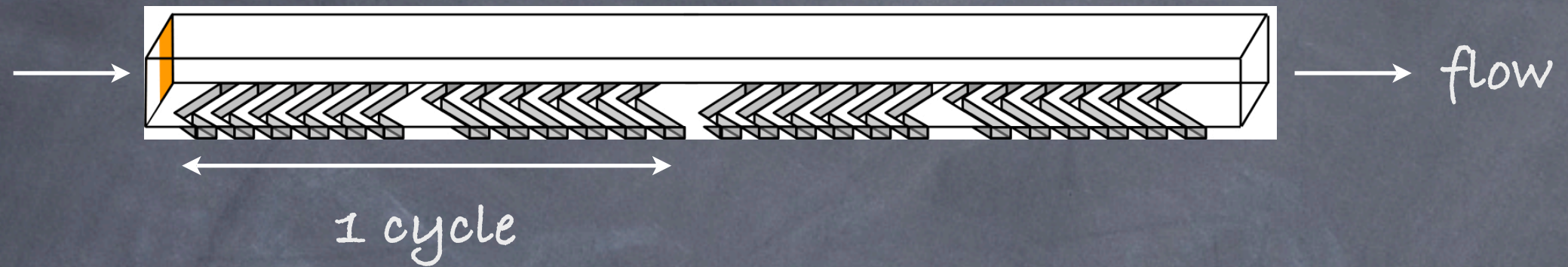
Somewhat more randomness

$$Re = 50$$



Smother fit

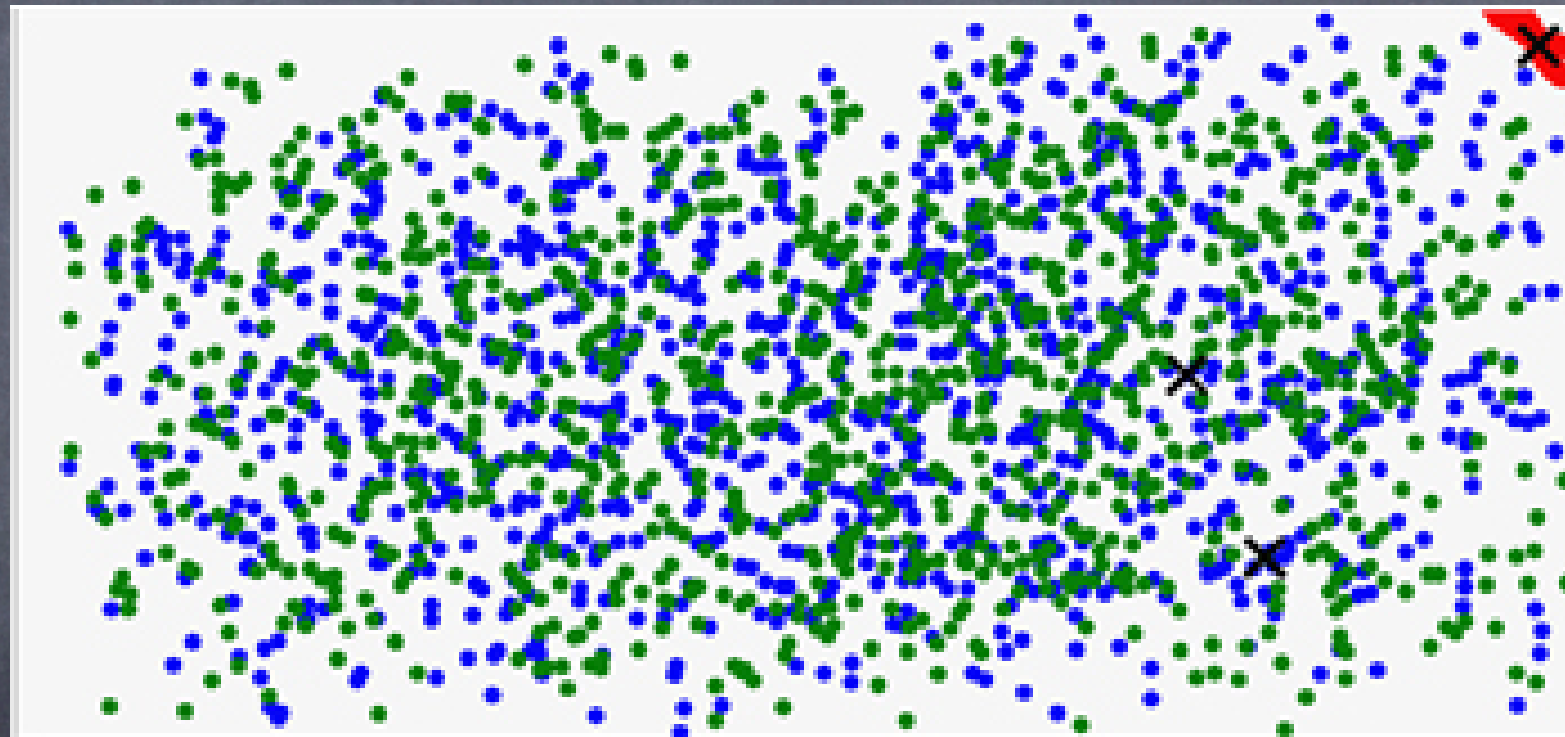
The 'herringbone' channel (Stroock & al. Science (2002))



100 μm

Poincaré section of initially segregated colored puffs

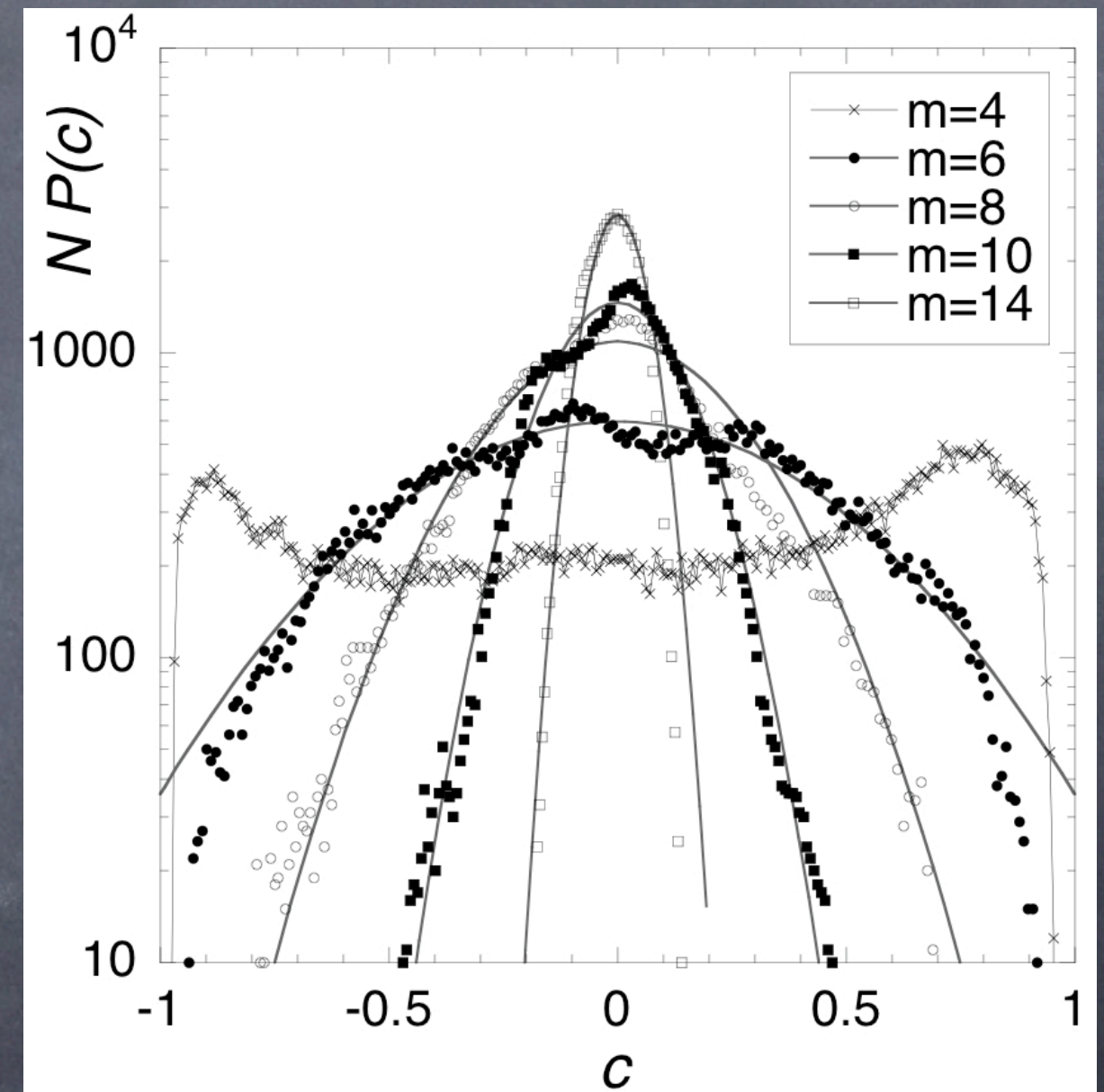
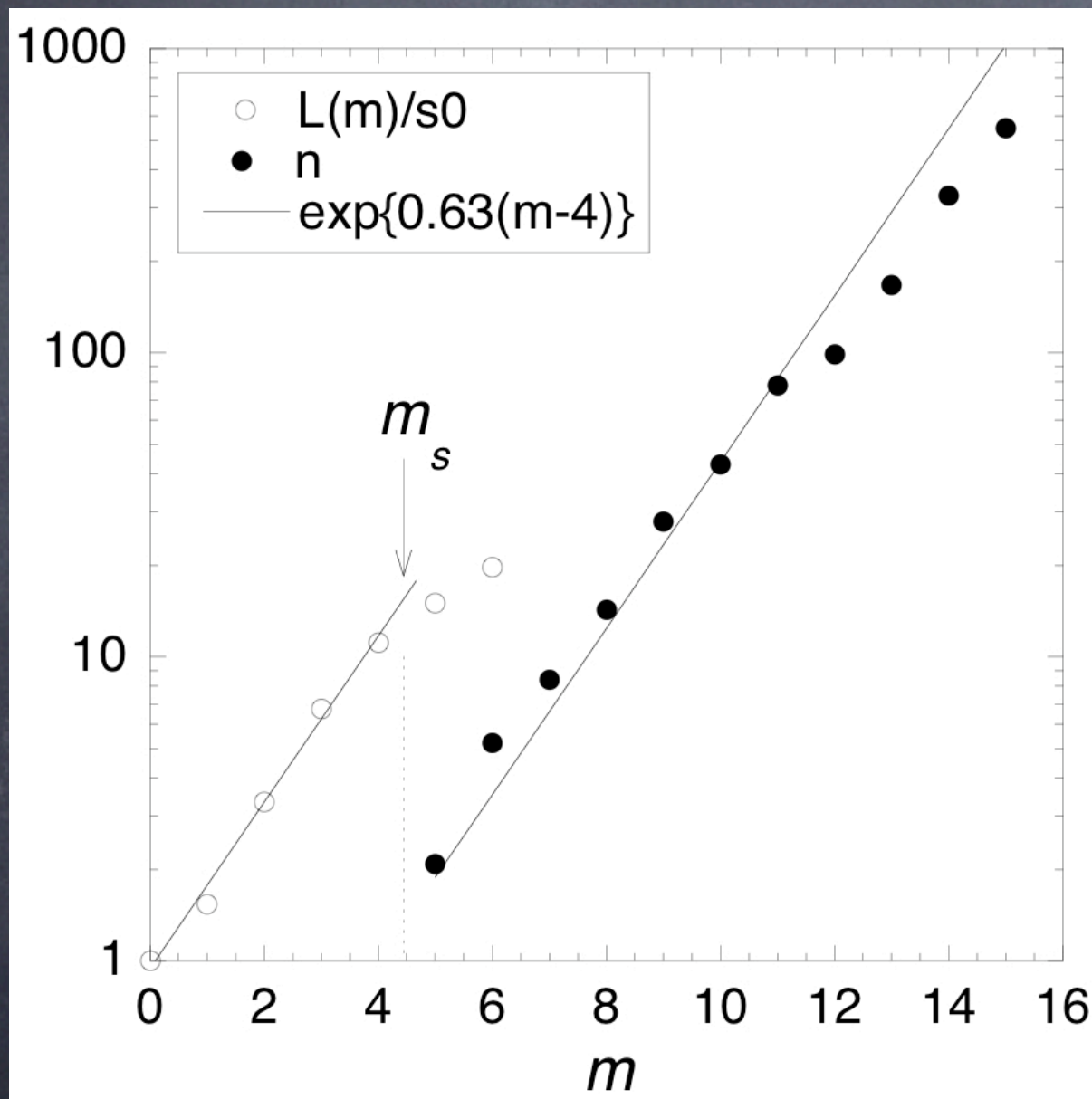
chaos synonym of random interactions



Stroock & McGraw, Proc. Roy. Soc. (2004)

convolutions about the mean: $c = C - \langle C \rangle$
 $\langle C \rangle = 1/2$

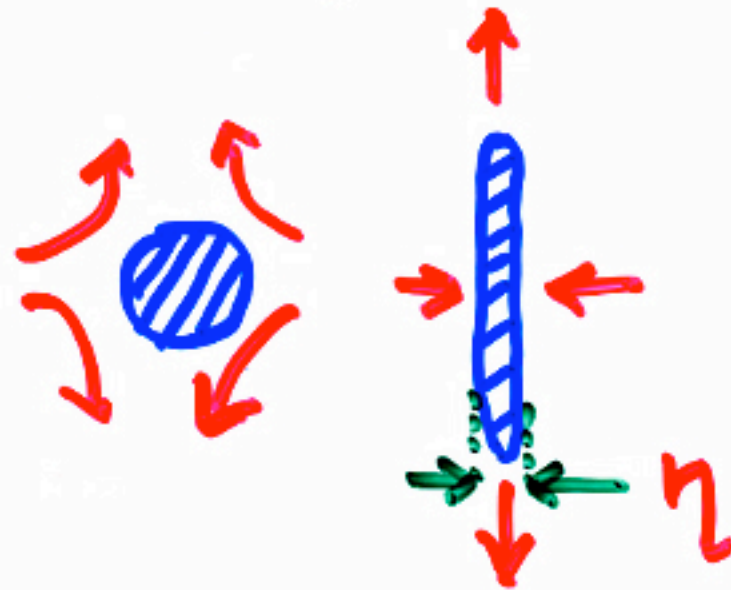
$P(c)$ = Bessel functions



Villermanx et al., PRL (2007)

A consequence

Equilibrium scales



$$\frac{\eta^2}{D} \sim \frac{r}{\delta u(r)}$$

- if $\delta u(r) \sim \gamma r$

$$\eta \sim \sqrt{\frac{D}{\gamma}}$$

Batchelor

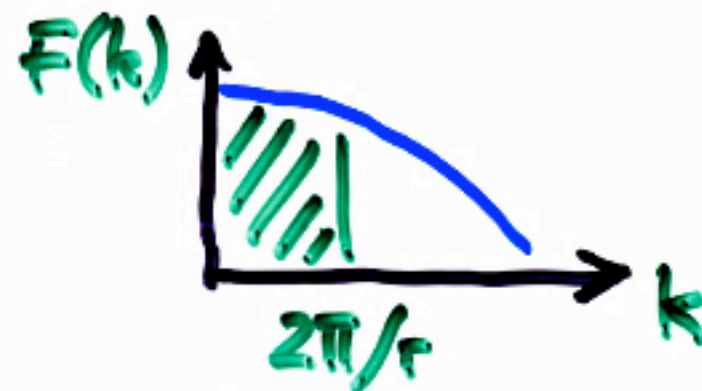
- if $\delta u(r) \sim (\epsilon r)^{1/3}$

$$\eta \sim \left(\frac{D^3}{\epsilon} \right)^{1/4}$$

Corrsin -
Obukhov

Coarse grained field

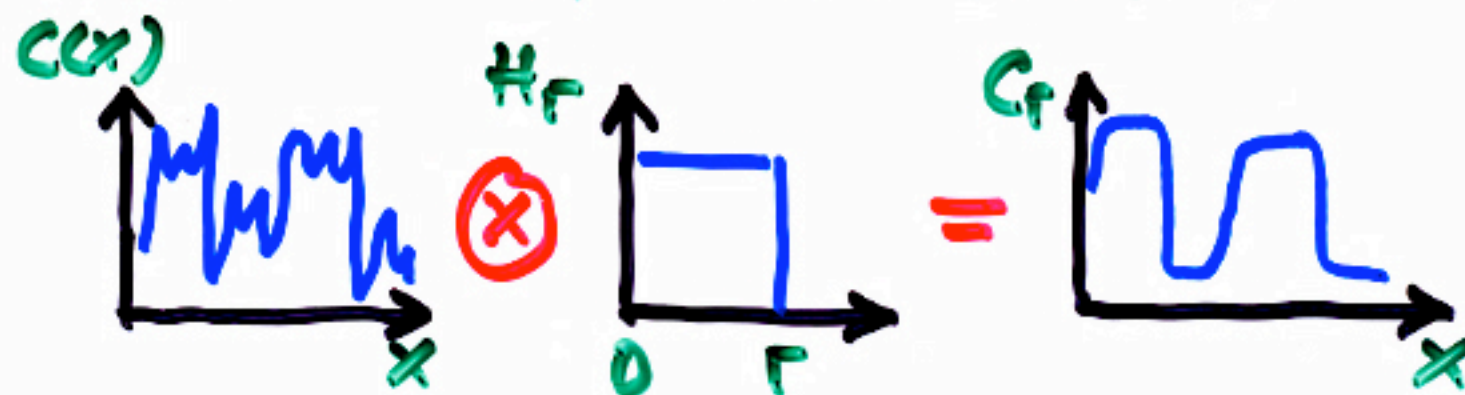
Spectrum: $|\int C(r) e^{ikr} dr|^2 = F(k)$

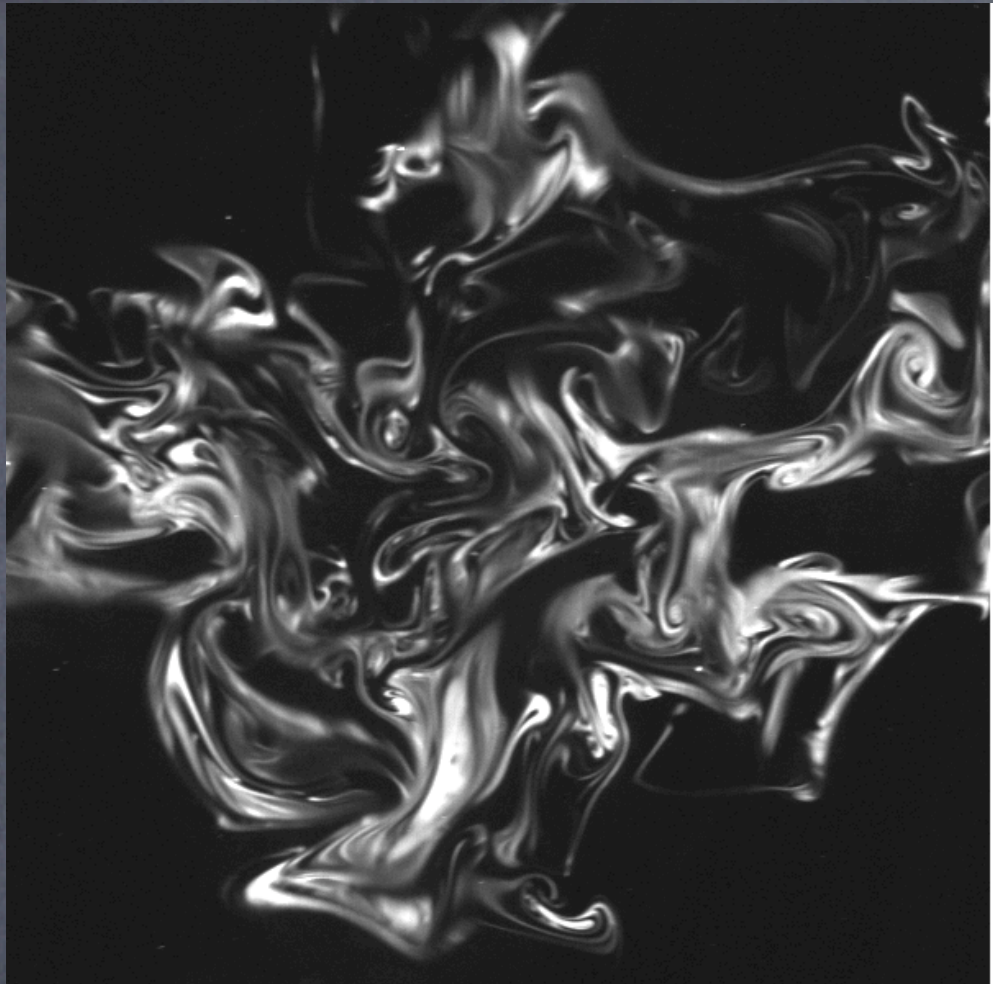


Coarse grained variance:

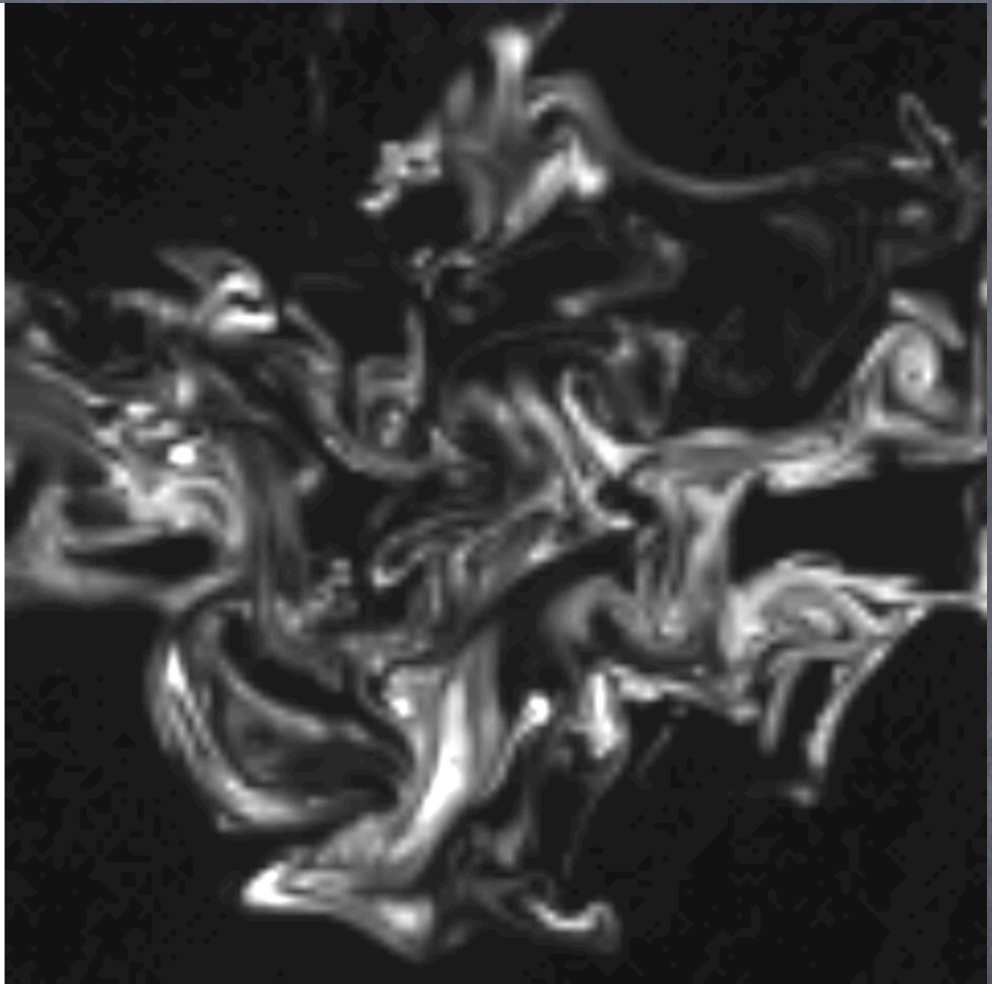
$$V_2(r) = \langle (C_r - \langle C_r \rangle)^2 \rangle$$

where $C_r = C(x) \otimes H_r(x)$



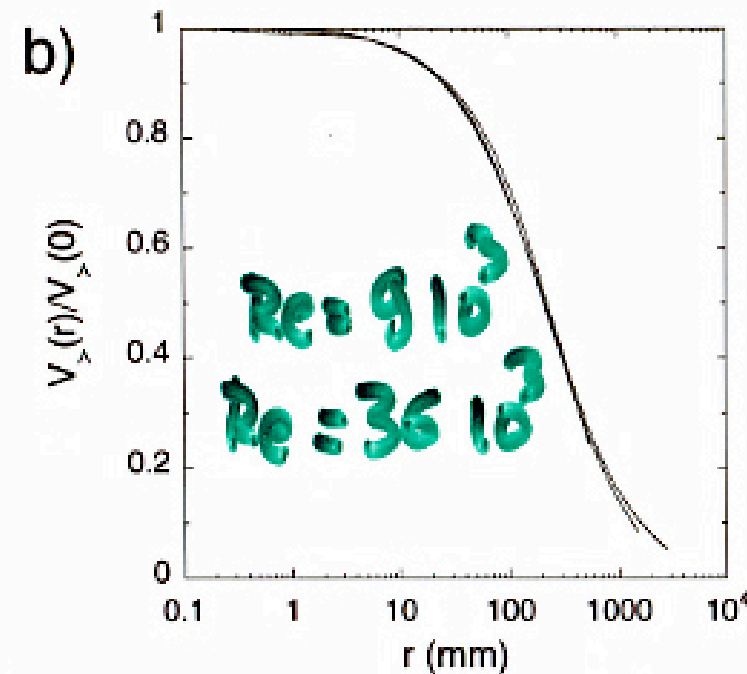
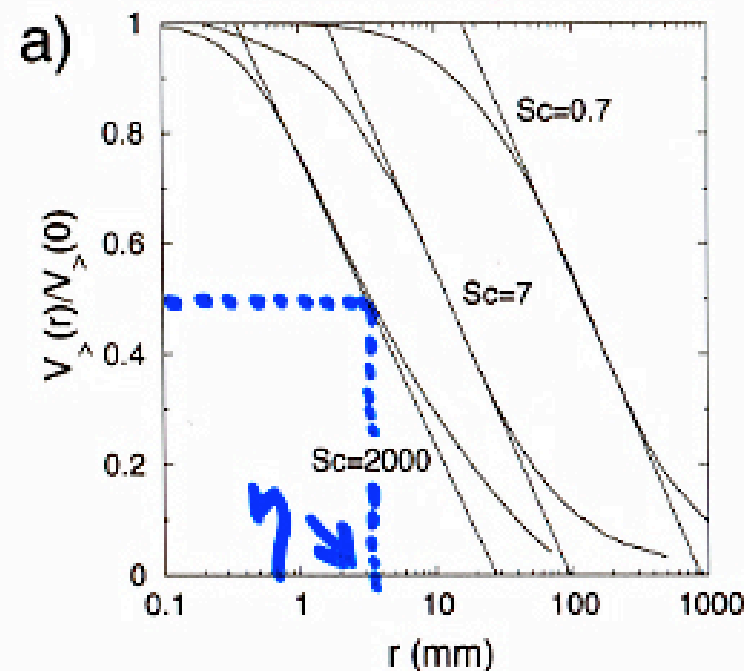


original

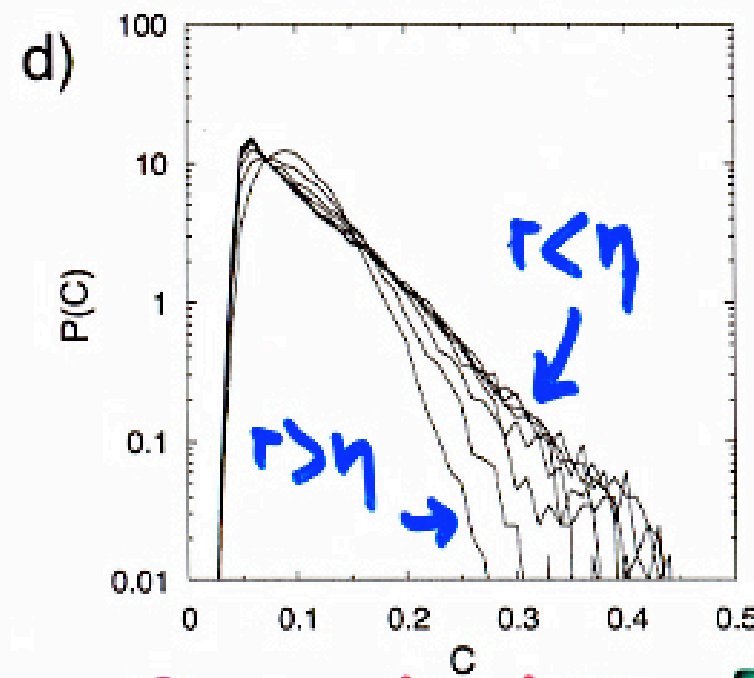
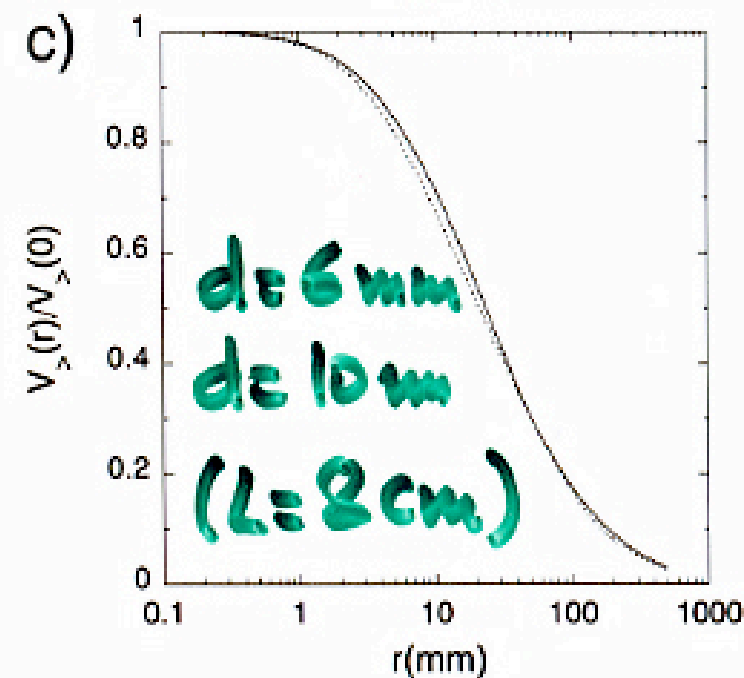


coarsened

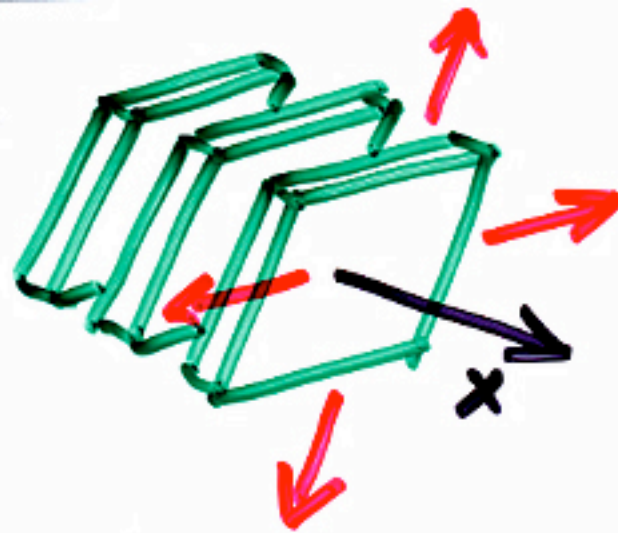
Variance $V_2(r)$



$V_2(r)$ depends
on Sc a)
is independent
of
 Re b)
Source size c)



Concentration Pdf



2 directions
of elongation

initially: $C(x, t=0) = 1 + \cos(k_0 x)$

then $\partial_t C - \partial_t \ln k(t) \cdot \partial_x C = D \partial_x^2 C$

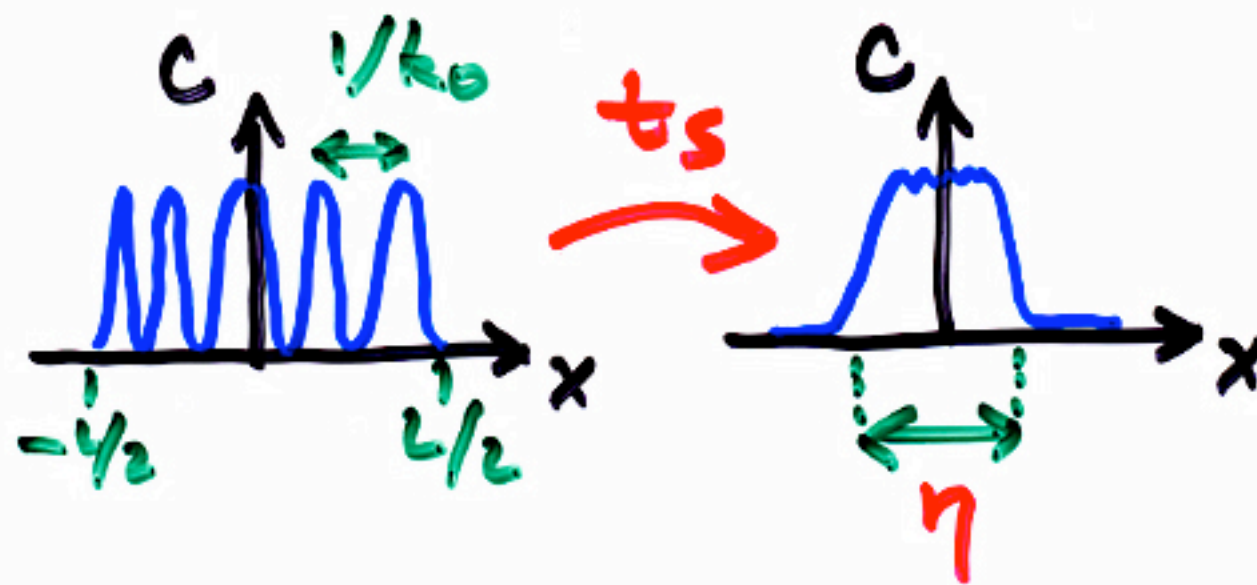
with $k(t) = k_0 (1 + \sigma t)^2$

$$\tau = \int dt' k(t')^2 = D k_0^2 \left(1 + \frac{2}{3} (\sigma t)^2 + \frac{1}{5} (\sigma t)^4 \right)$$

$$\xi = x k(t)$$

$$\partial_t C = \partial_{\xi^2} C$$

pure diffusion
equation



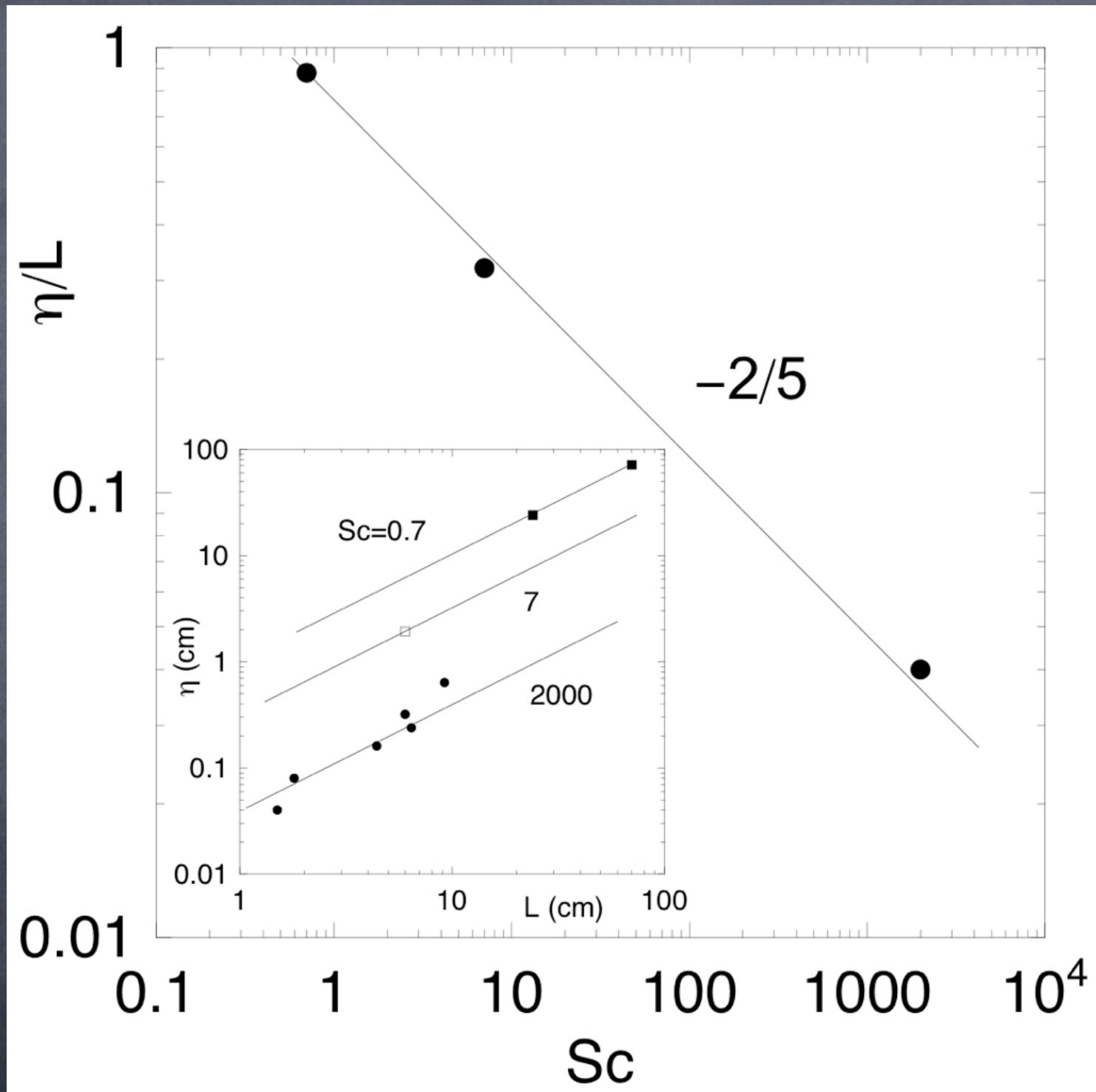
Mixing time: $t_s = \frac{1}{\sigma} Pe^{1/5}$; $Pe = \frac{U}{D k_0^2} \gg 1$

Taylor scale: $k_0 = \sqrt{\frac{U}{2\nu}}$

$\eta = \frac{L}{(1 + \sigma/5)^{2/5}} \approx L Sc^{-2/5}$

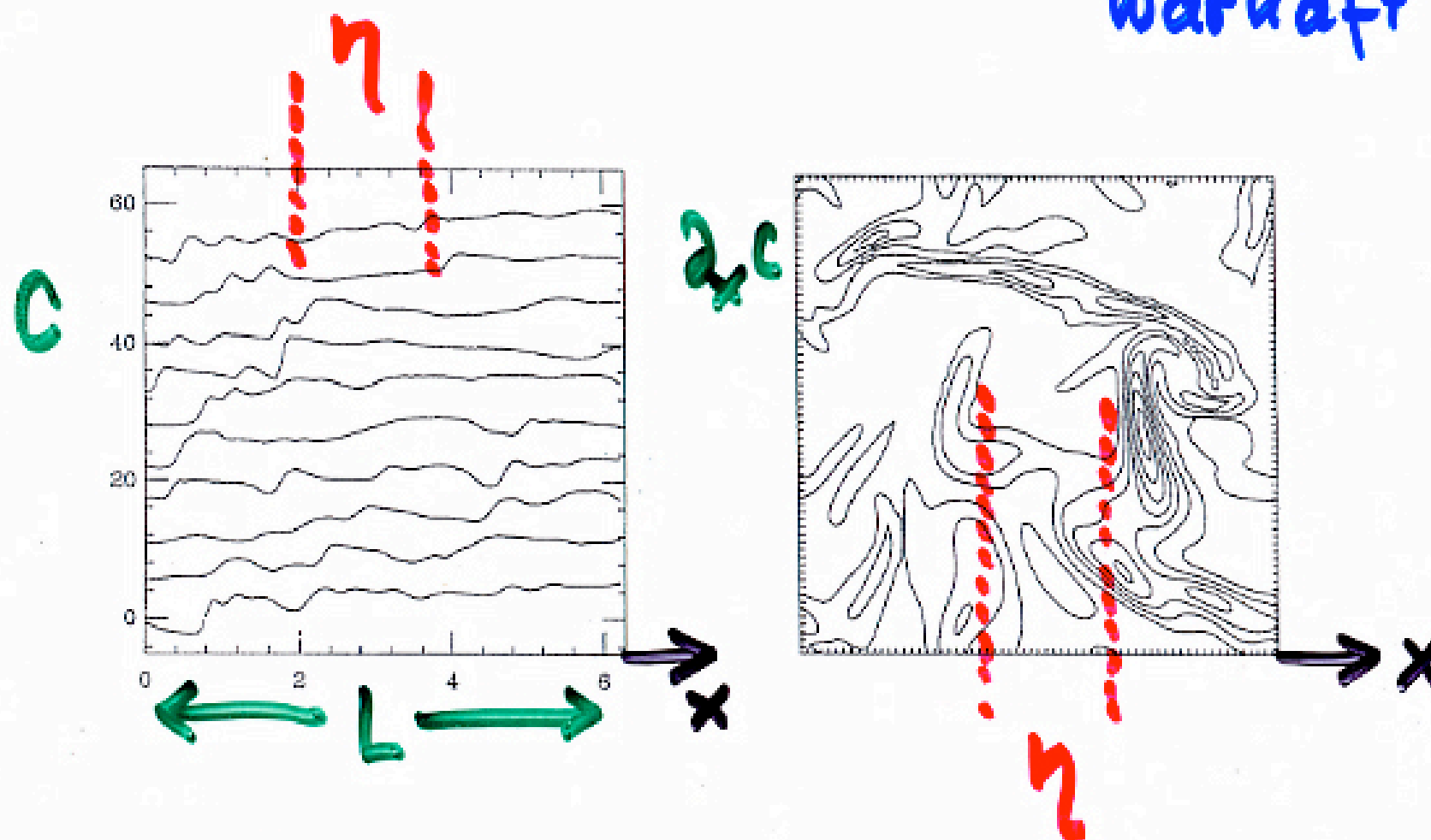
Coarse grained scale

Natural coarse grained
scale $\eta = L Sc^{-2/5}$



Villiermaux & Duplat PRL 97, 144506 (2006)

Plateaux, in shear flows
Pumir 194
Warhaft '00



Mixing is not only an academic topic:
Search for the source !

('Infotaxis', Vergassola & al., Nature 2007)

